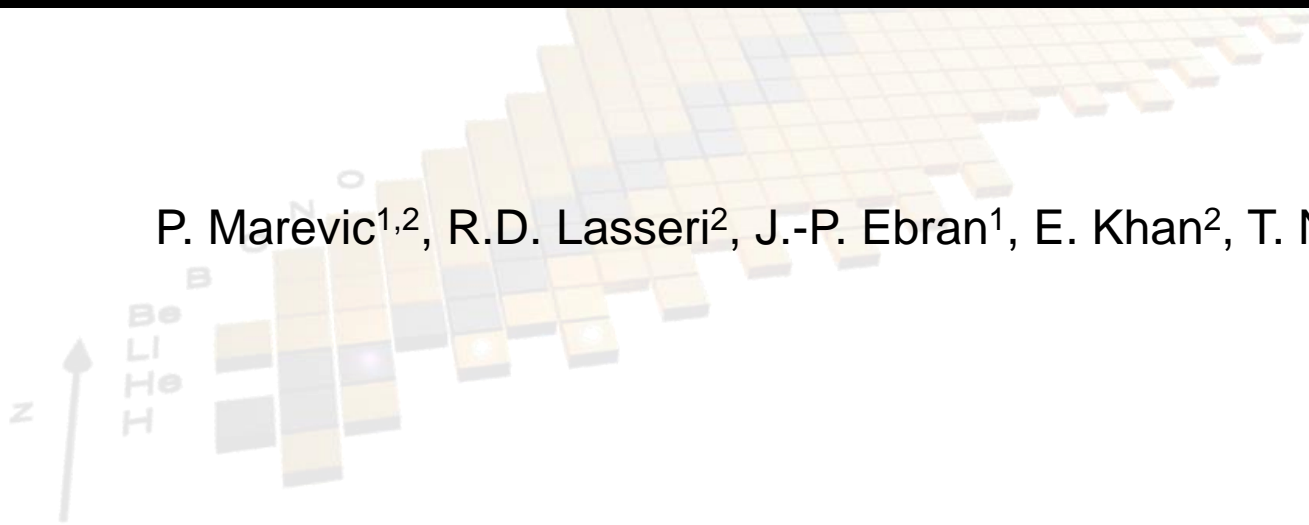




# -Nuclear Clustering in the Energy Density Functional Approach-



P. Marevic<sup>1,2</sup>, R.D. Lasserri<sup>2</sup>, J.-P. Ebran<sup>1</sup>, E. Khan<sup>2</sup>, T. Niksic<sup>3</sup>, D.Vretenar<sup>3</sup>

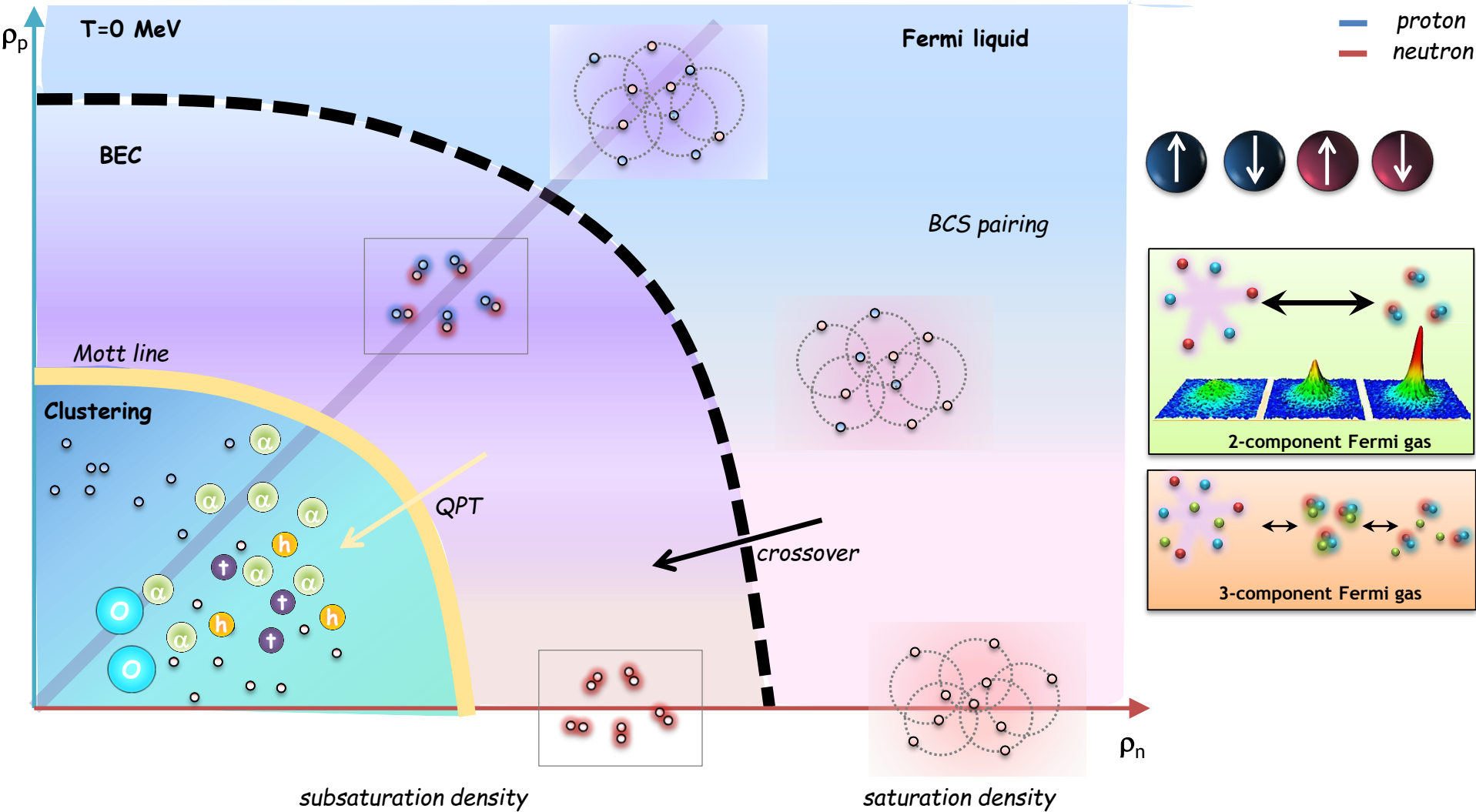


<sup>1</sup> CEA,DAM,DIF

<sup>2</sup> IPNO

<sup>3</sup> University of Zagreb

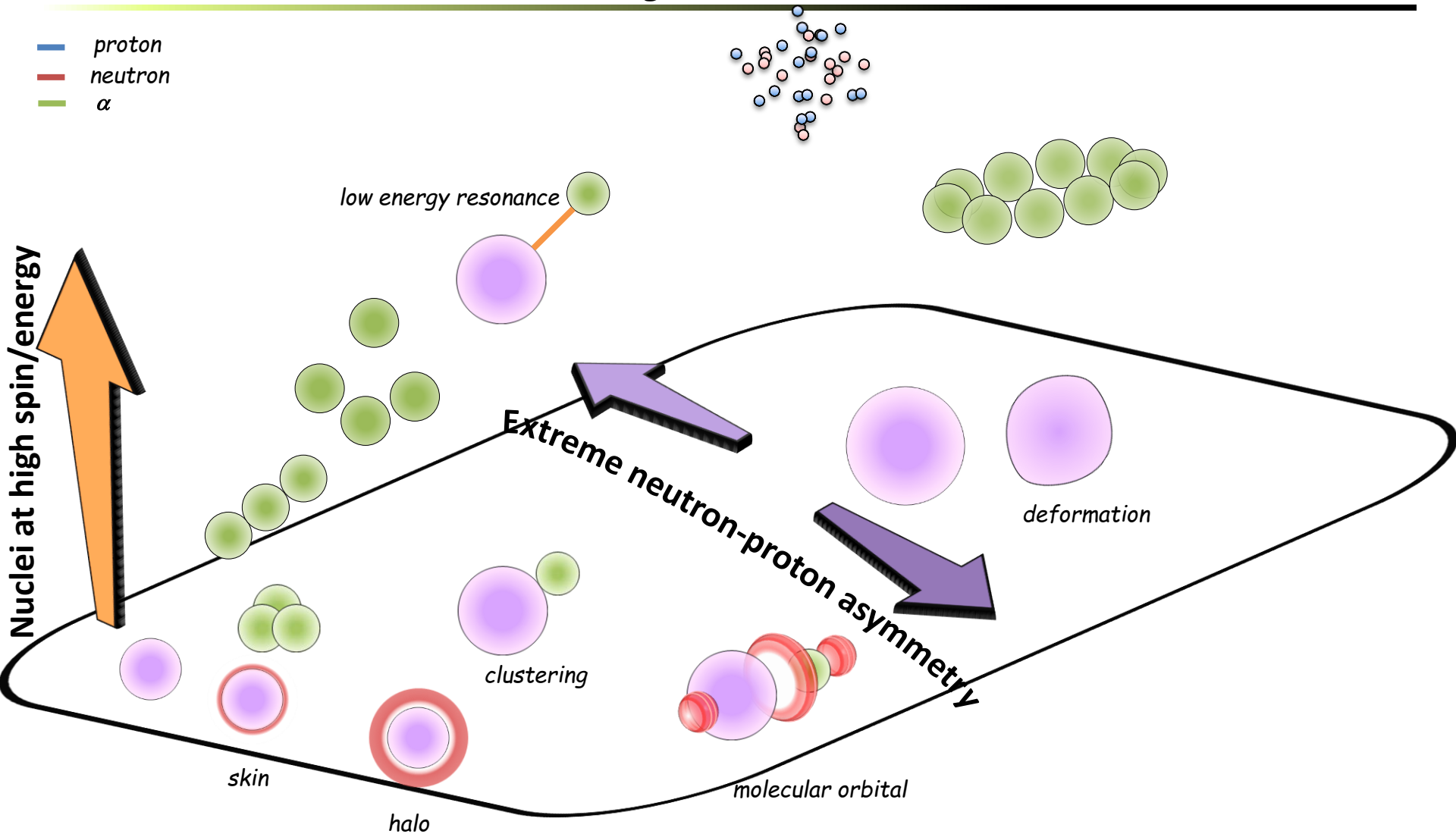
# Nuclear systems as a mixture of 4 types of fermion



⇒ These features leave fingerprints in finite nuclei

# Emergence of ordered sub-structures inside nuclei

- proton
- neutron
- $\alpha$



⇒ Nuclear EDFs provide a unified and consistent description of these various properties

**1** Theoretical framework : EDF

**2** Pair correlations

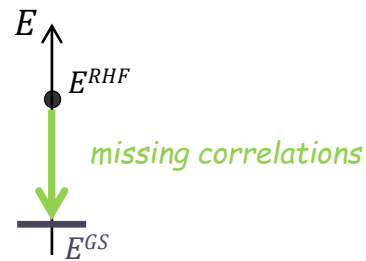
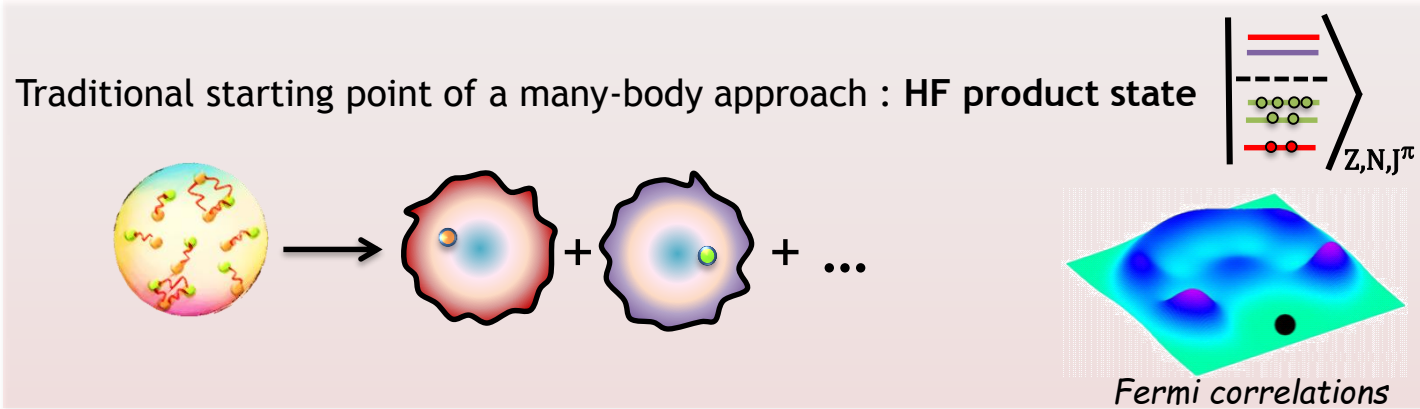
**3** Nuclear clustering



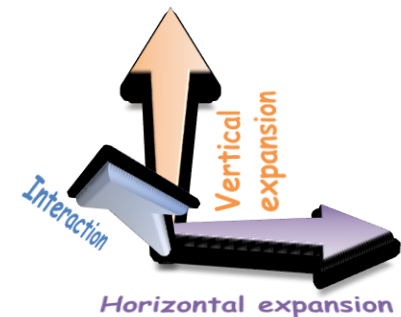
- ① *Theoretical framework : EDF-*

# ⊗ Nuclear many-body problem : strategies

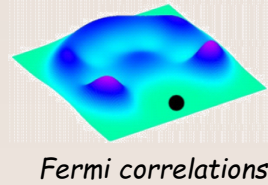
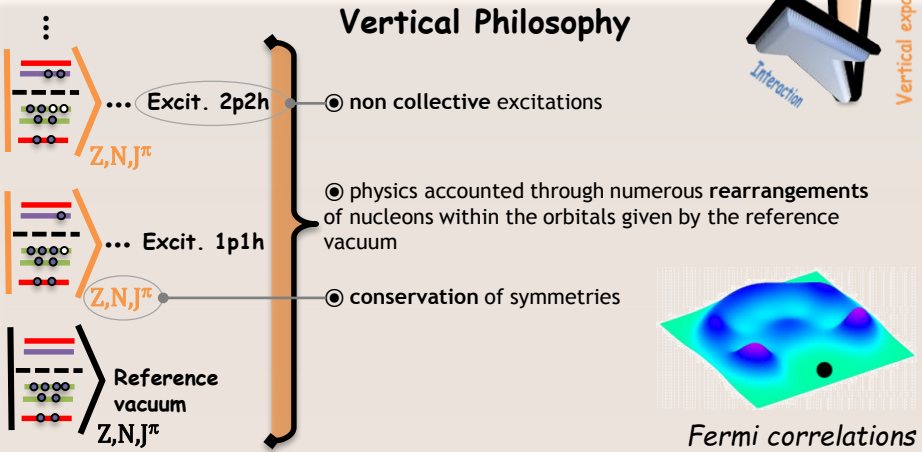
⇒ Many-body approaches as implementation of different strategies to apprehend the many-body problem



⇒ How to incorporate such correlations starting with a HF product state ?



# Nuclear many-body problem : strategies



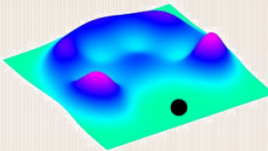
$$|Z, N, J^\pi\rangle = e_0 |Z, N, J^\pi\rangle + e_i^{1p-1h} |Z, N, J^\pi\rangle + e_j^{2p-2h} |Z, N, J^\pi\rangle + \dots$$

The equation shows the ground state  $|Z, N, J^\pi\rangle$  (represented by a globe icon) as a sum of components: the reference vacuum  $e_0 |Z, N, J^\pi\rangle$ , a 1p-1h excitation  $e_i^{1p-1h} |Z, N, J^\pi\rangle$ , and a 2p-2h excitation  $e_j^{2p-2h} |Z, N, J^\pi\rangle$ , followed by higher-order terms.

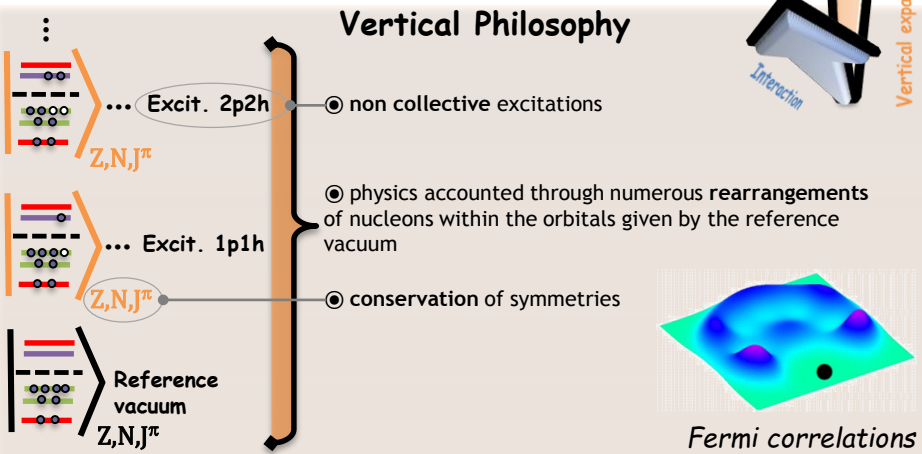
Major contribution

Dynamical correlations

- in all nuclei
- $|HF\rangle \sim |globe\rangle$



# Nuclear many-body problem : strategies

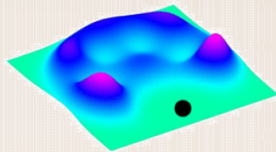


$$|Z, N, J^\pi\rangle = e_0 |Z, N, J^\pi\rangle + e_i^{1p-1h} |Z, N, J^\pi\rangle + e_j^{2p-2h} |Z, N, J^\pi\rangle + \dots$$

comparable contributions

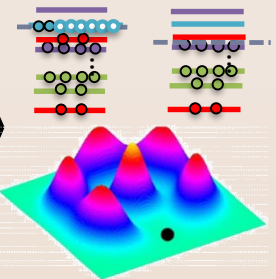
Dynamical correlations

- in all nuclei
- $|HF\rangle \sim |\text{globe}\rangle$



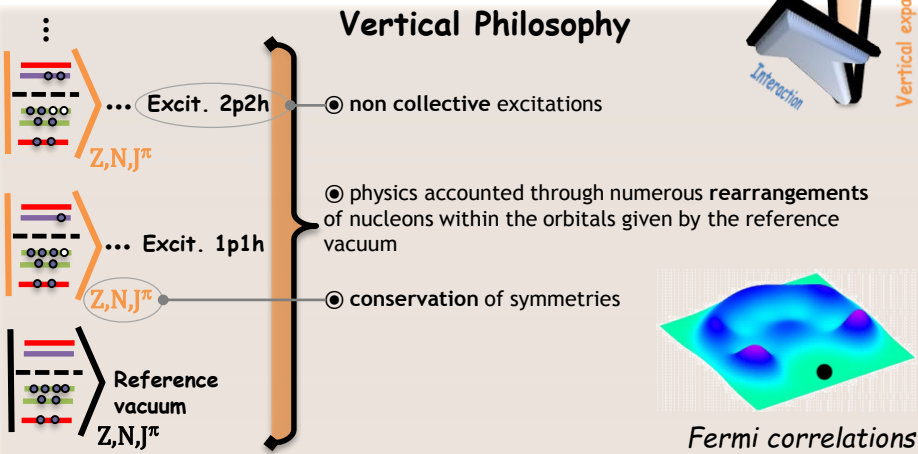
Nondynamical correlations

- open shell systems
- $|HF\rangle \neq |\text{globe}\rangle$

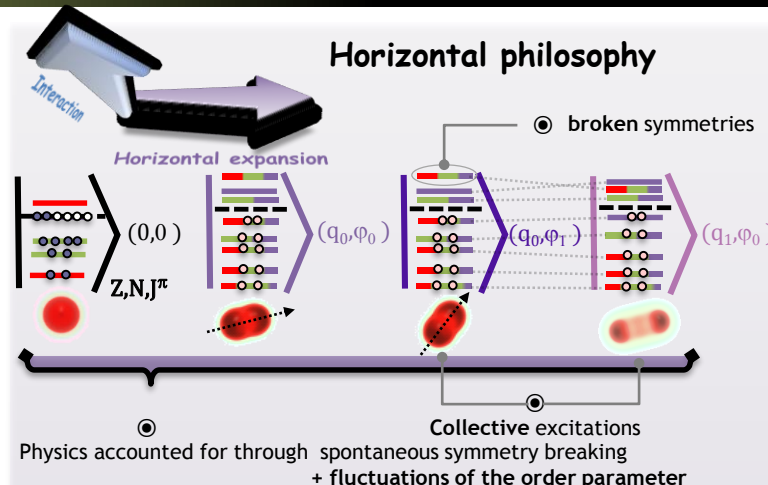
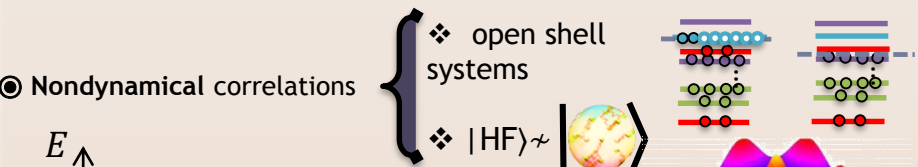
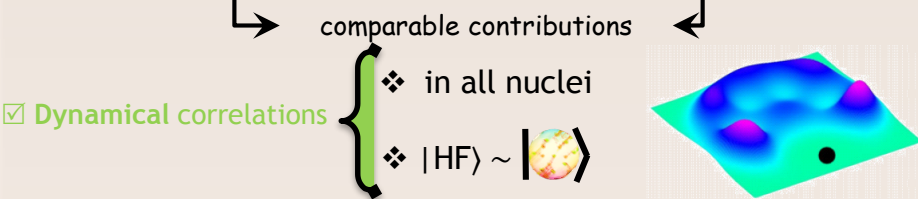




# Nuclear many-body problem : strategies

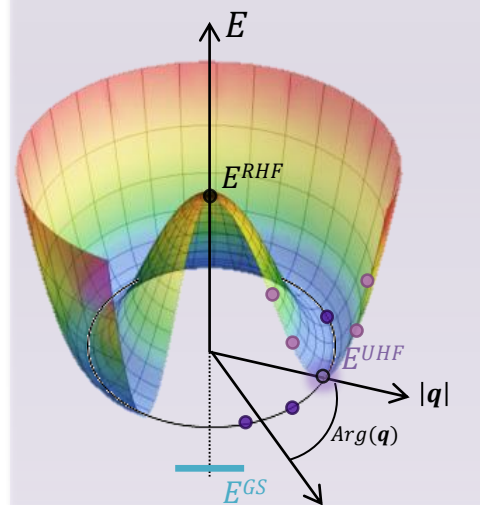


$$| \text{Nucleus} \rangle_{Z, N, J^\pi} = e_0 | \text{Reference vacuum} \rangle_{Z, N, J^\pi} + e_i^{1p-1h} | \text{Excit. 1p1h} \rangle_{Z, N, J^\pi} + e_j^{2p-2h} | \text{Excit. 2p2h} \rangle_{Z, N, J^\pi} + \dots$$



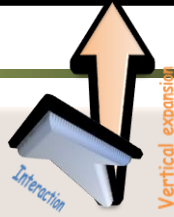
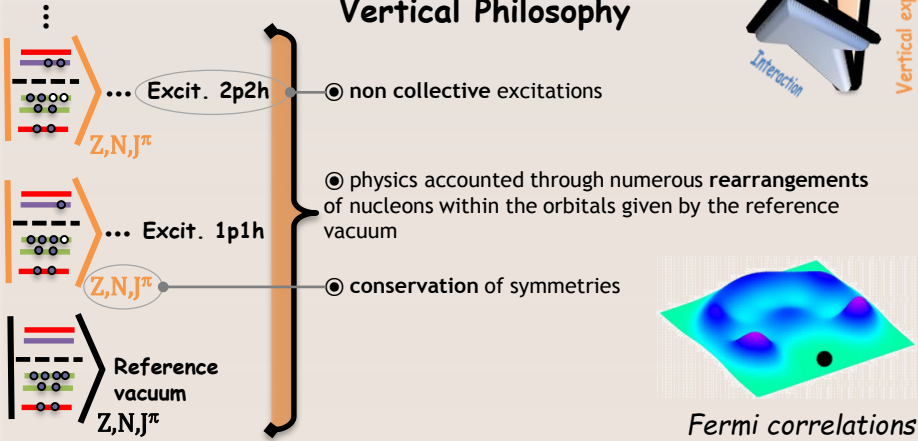
dynamical correlations transferred in  $H_{eff}(q)$

$$| \text{Nucleus} \rangle = \int dq f(q) | \text{Excit. } (q) \rangle (|q|, \text{arg}(q))$$

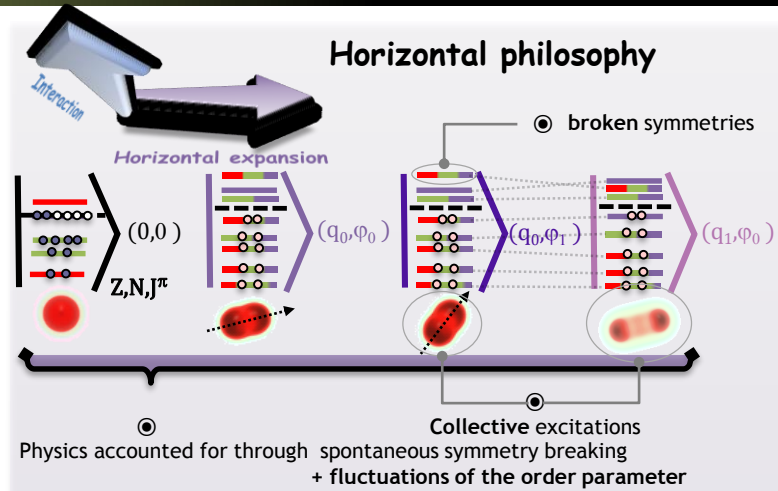


# Nuclear many-body problem : strategies

## Vertical Philosophy



## Horizontal philosophy



dynamical correlations transferred in  $H_{eff}(q)$

$$|\text{Nucleus}\rangle = \int dq f(q) | |q|, \arg(q) \rangle$$

$$|\text{Nucleus}\rangle = e_0 |Z, N, J^\pi\rangle + e_i^{1p-1h} |Z, N, J^\pi\rangle + e_i^{2p-2h} |Z, N, J^\pi\rangle + \dots$$

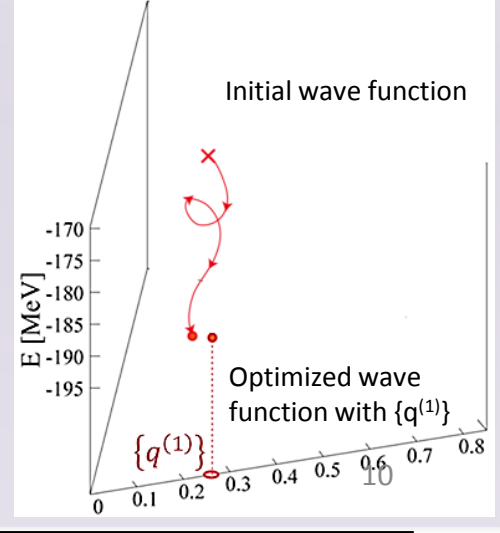
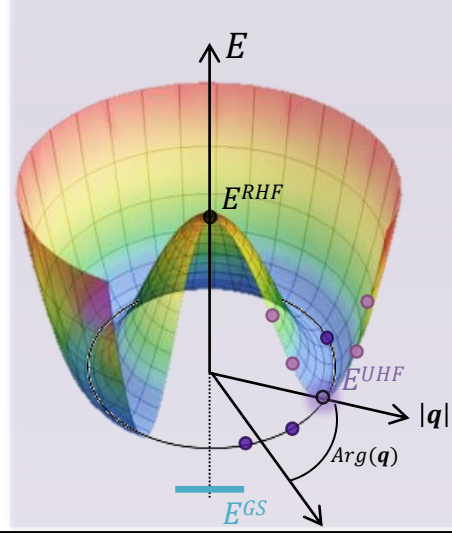
comparable contributions

### Dynamical correlations

- in all nuclei
- $|\text{HF}\rangle \sim |\text{Nucleus}\rangle$

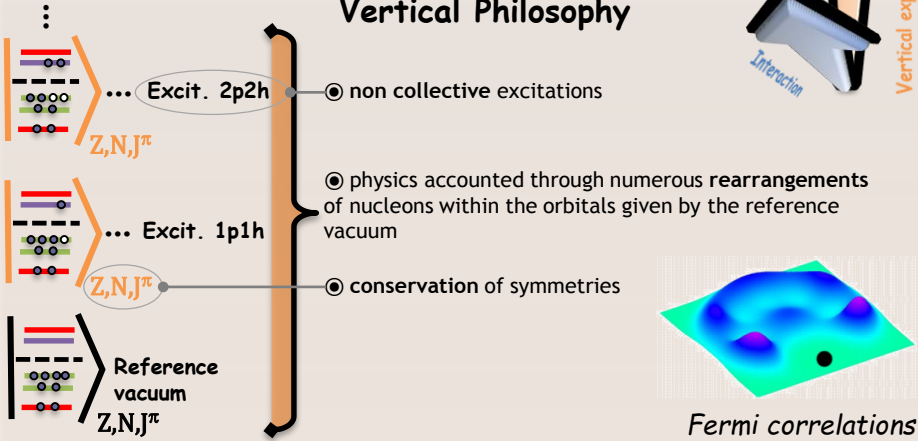
### Nondynamical correlations

- open shell systems
- $|\text{HF}\rangle \not\sim |\text{Nucleus}\rangle$



# Nuclear many-body problem : strategies

## Vertical Philosophy



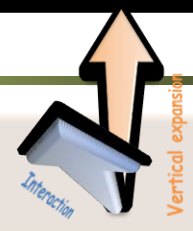
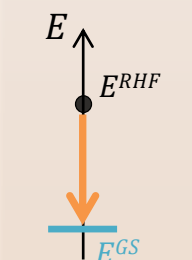
$$| \text{Nucleus} \rangle_{Z, N, J^\pi} = e_0 | \text{vacuum} \rangle_{Z, N, J^\pi} + e_i^{1p-1h} | \text{excit. 1p1h} \rangle_{Z, N, J^\pi} + e_j^{2p-2h} | \text{excit. 2p2h} \rangle_{Z, N, J^\pi} + \dots$$

comparable contributions

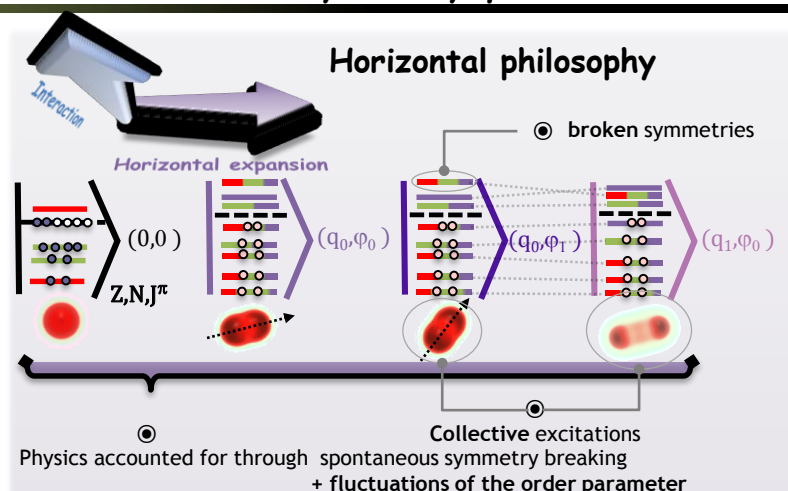
- in all nuclei
- $| \text{HF} \rangle \sim | \text{Nucleus} \rangle$

Nondynamical correlations

- open shell systems
- $| \text{HF} \rangle \not\sim | \text{Nucleus} \rangle$

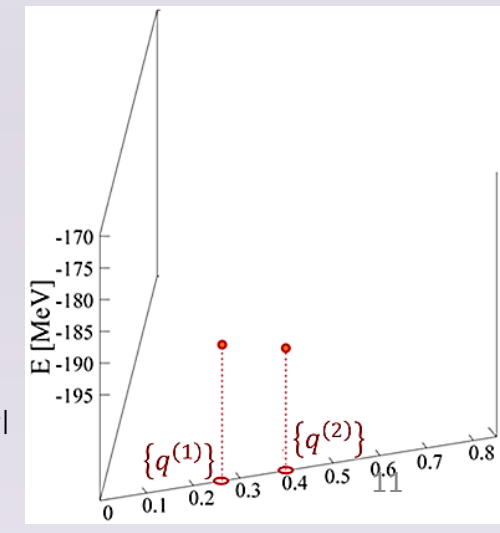
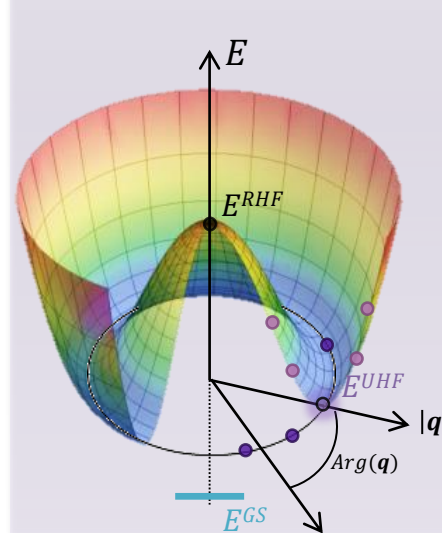


## Horizontal philosophy



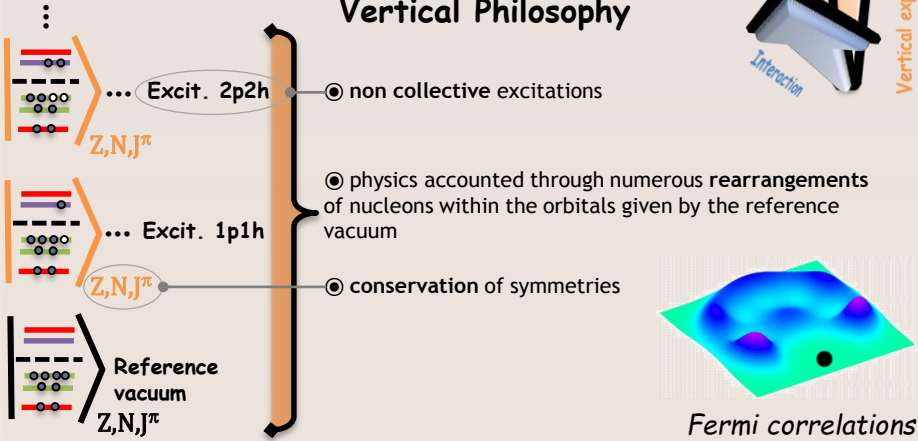
dynamical correlations transferred in  $H_{\text{eff}}(q)$

$$| \text{Nucleus} \rangle = \int dq f(q) | (|q|, \text{arg}(q)) \rangle$$



# Nuclear many-body problem : strategies

## Vertical Philosophy



$$| \text{Nucleus} \rangle_{Z, N, J^\pi} = e_0 | \text{Reference vacuum} \rangle_{Z, N, J^\pi} + e_i^{1p-1h} | \text{Excit. 1p1h} \rangle_{Z, N, J^\pi} + e_j^{2p-2h} | \text{Excit. 2p2h} \rangle_{Z, N, J^\pi} + \dots$$

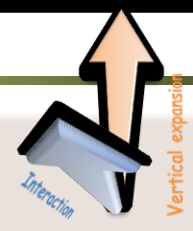
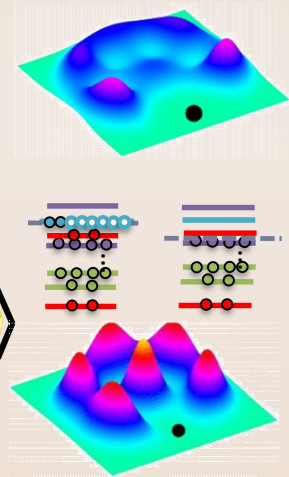
comparable contributions

### Dynamical correlations

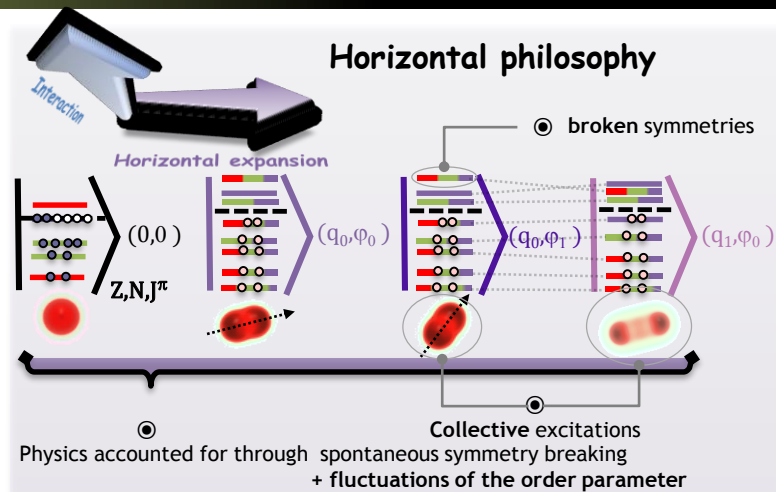
- in all nuclei
- $| \text{HF} \rangle \sim | \text{Nucleus} \rangle$

### Nondynamical correlations

- open shell systems
- $| \text{HF} \rangle \not\sim | \text{Nucleus} \rangle$

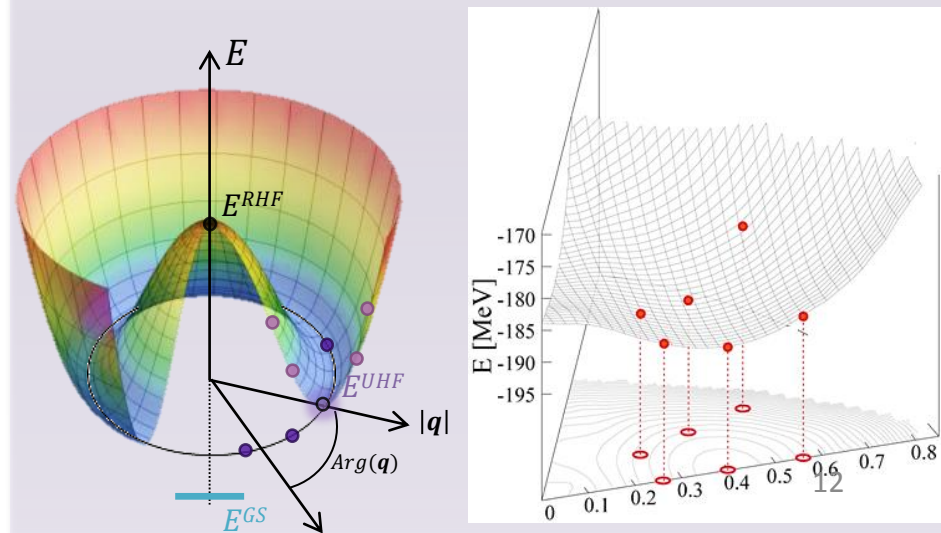


## Horizontal philosophy



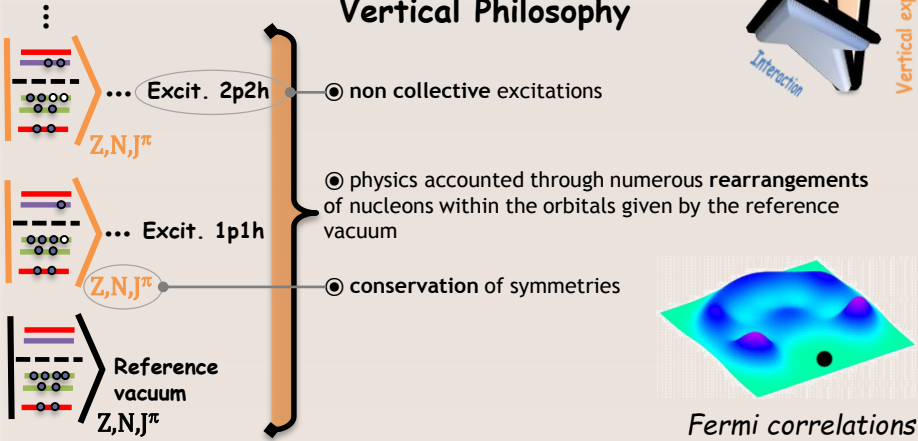
dynamical correlations transferred in  $H_{\text{eff}}(q)$

$$| \text{Nucleus} \rangle = \int dq f(q) | (|q|, \text{arg}(q)) \rangle$$



# Nuclear many-body problem : strategies

## Vertical Philosophy



$$| \text{Nucleus} \rangle_{Z, N, J^\pi} = e_0 | \text{Reference vacuum} \rangle_{Z, N, J^\pi} + e_i^{1p-1h} | \text{Excit. 1p1h} \rangle_{Z, N, J^\pi} + e_j^{2p-2h} | \text{Excit. 2p2h} \rangle_{Z, N, J^\pi} + \dots$$

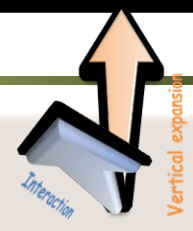
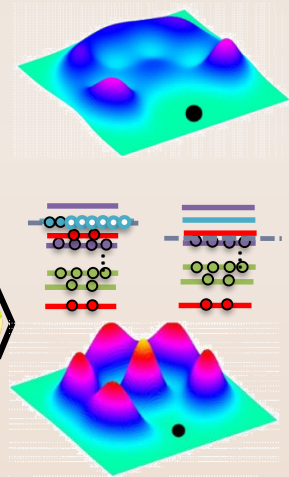
comparable contributions

### Dynamical correlations

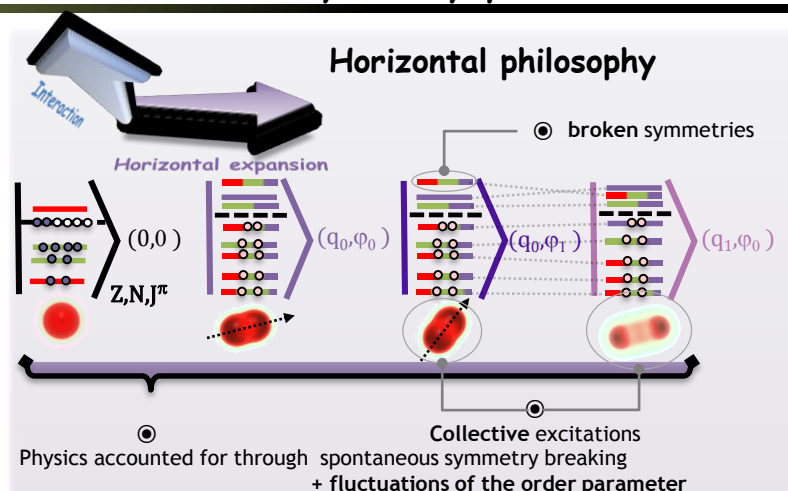
- in all nuclei
- $| \text{HF} \rangle \sim | \text{Nucleus} \rangle$

### Nondynamical correlations

- open shell systems
- $| \text{HF} \rangle \not\sim | \text{Nucleus} \rangle$

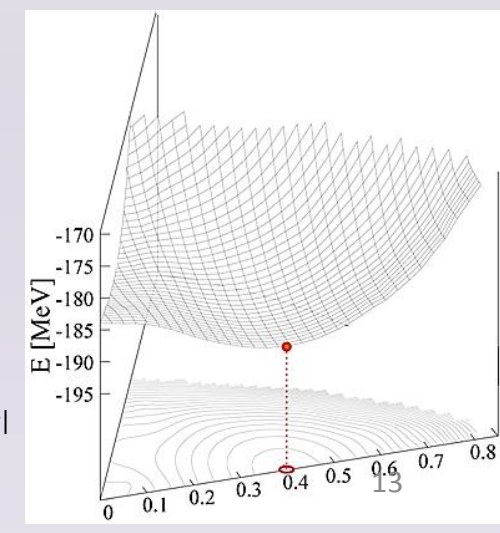
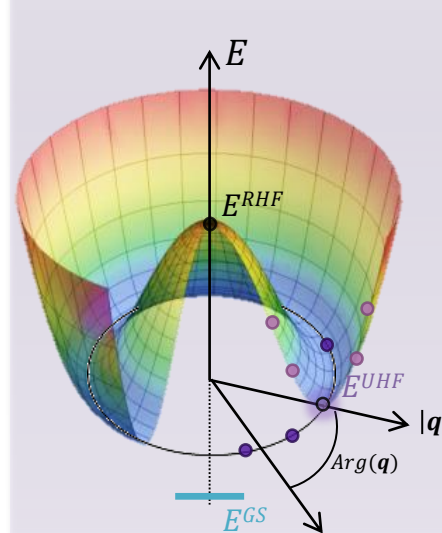


## Horizontal philosophy



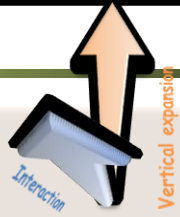
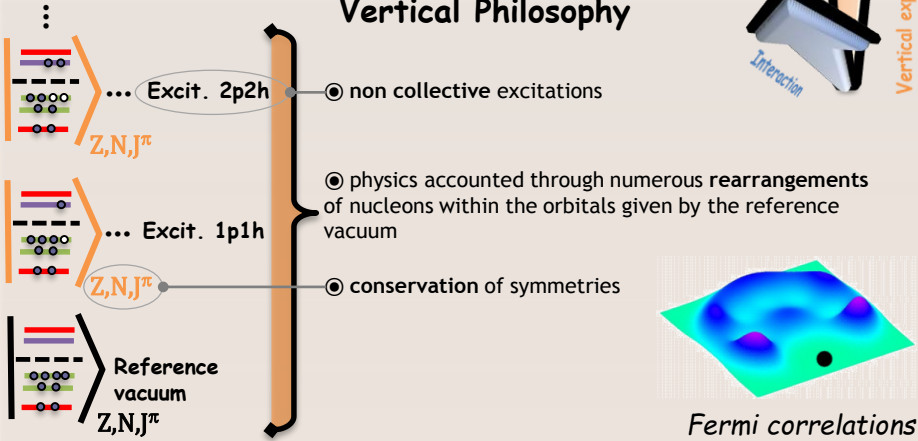
dynamical correlations transferred in  $H_{eff}(q, \phi)$

$$| \text{Nucleus} \rangle = \int dq f(q) | (|q|, \arg(q)) \rangle$$

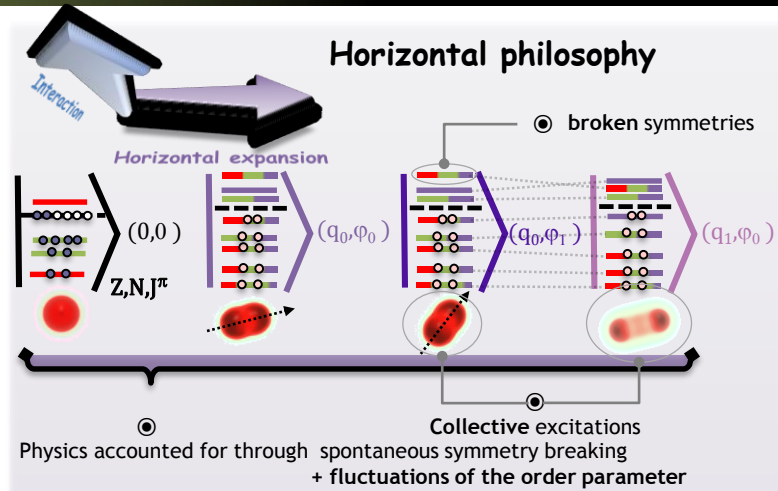


# Nuclear many-body problem : strategies

## Vertical Philosophy



## Horizontal philosophy



$$|Z, N, J^\pi\rangle = e_0 |Z, N, J^\pi\rangle + e_i^{1p-1h} |Z, N, J^\pi\rangle + e_j^{2p-2h} |Z, N, J^\pi\rangle + \dots$$

comparable contributions

### Dynamical correlations

- in all nuclei
- $|HF\rangle \sim |\text{globe}\rangle$

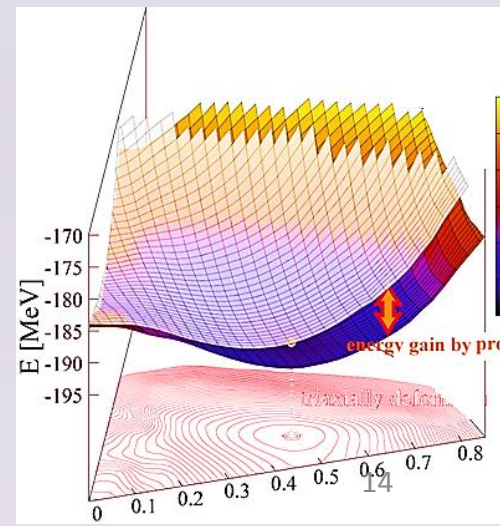
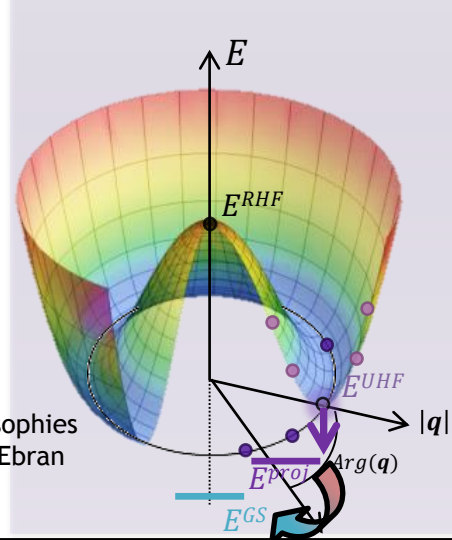
### Nondynamical correlations

- open shell systems
- $|HF\rangle \not\sim |\text{globe}\rangle$



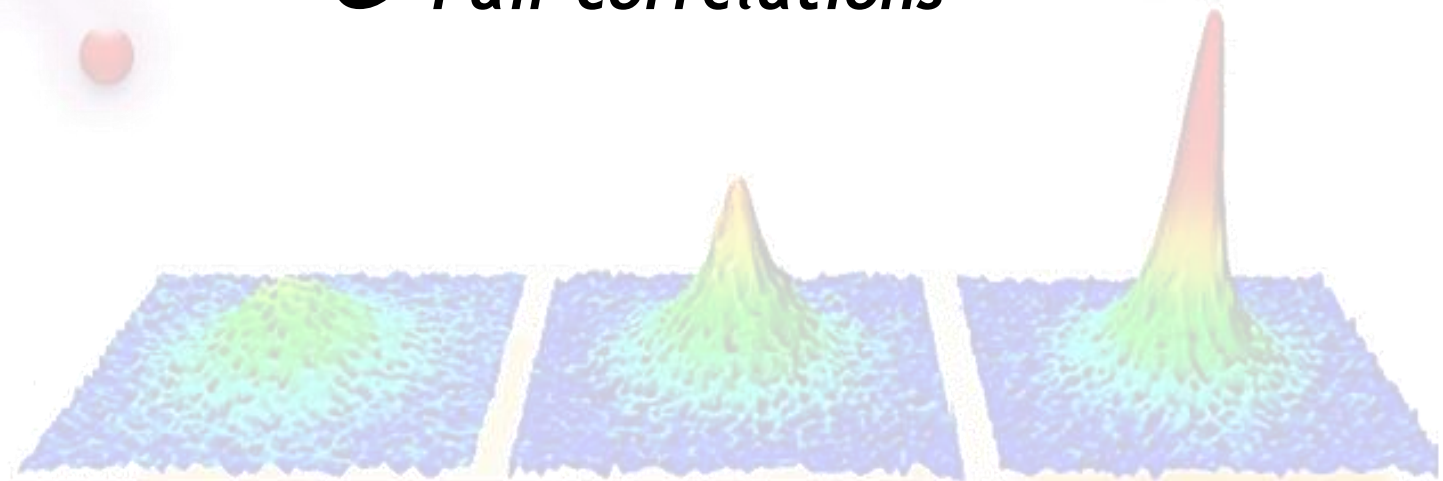
dynamical correlations transferred in  $H_{eff}(q)$

$$|\text{globe}\rangle = \int dq f(q) | |q\rangle (|q|, \arg(q))$$



Novel many-body approaches combining both philosophies  
 P. Arthuis, J. Ripoché, T. Duguet, D. Lacroix, J.-P. Ebran  
 J. Ripoché et al PRC 95, 014326 (2017)  
 T. Duguet et al EPJA 51, 162 (2015)

- **2** *Pair correlations* -



★ How to describe a pair of nucleons in the medium ?

Reduced density matrix

- All the information of a many-body system is contained in its Nth-order density matrix

$$D_N(1, \dots, N; 1', \dots, N') = \Psi(1, \dots, N) \Psi^*(1', \dots, N')$$

- Only keep information about p-“cluster” embedded in the medium composed by the other N-p particles :

$$\Gamma_p(1, \dots, p; 1', \dots, p') = \binom{N}{p} \int d(p+1) \dots dN \Psi(1, \dots, N) \Psi^*(1', \dots, p', p+1, \dots, N)$$

- 2-RDM : eigenfunctions provide an in-medium pair wave function

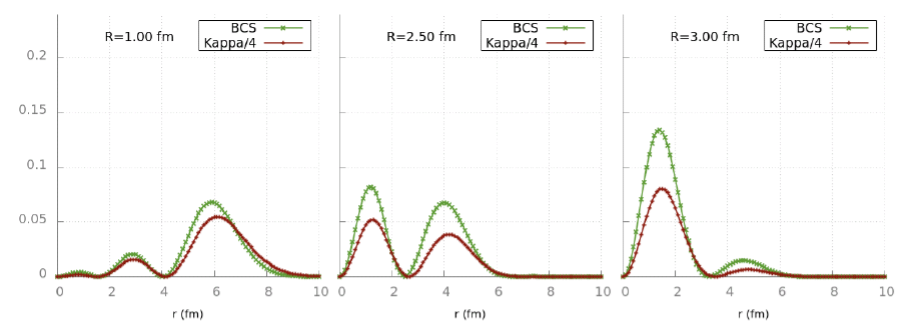
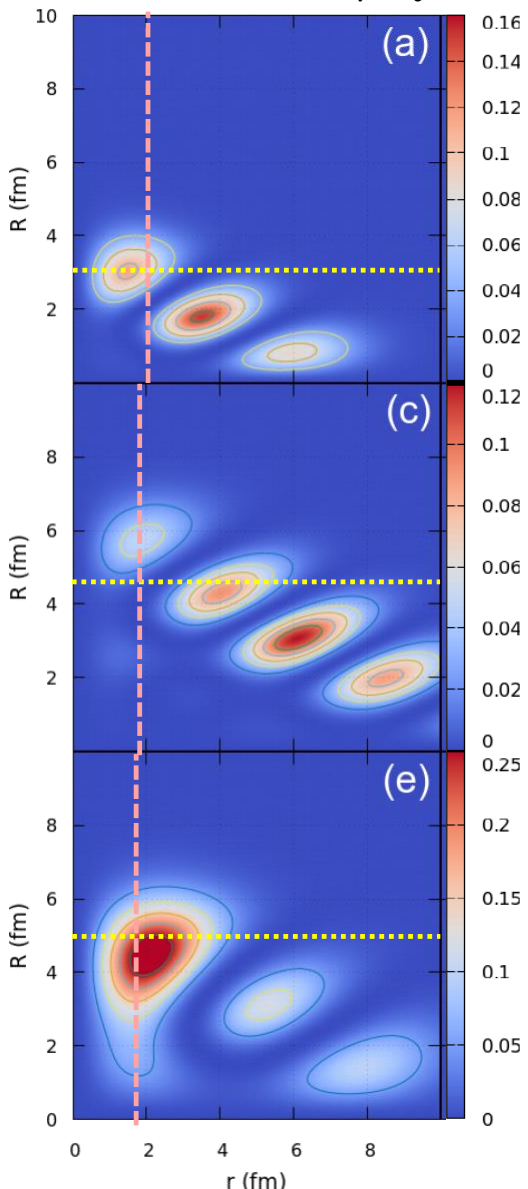
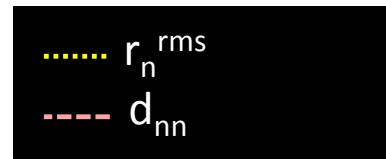
Form of the A-body wave function

- BCS assumption: all the pairs of fermions occupy the same pair wave function :

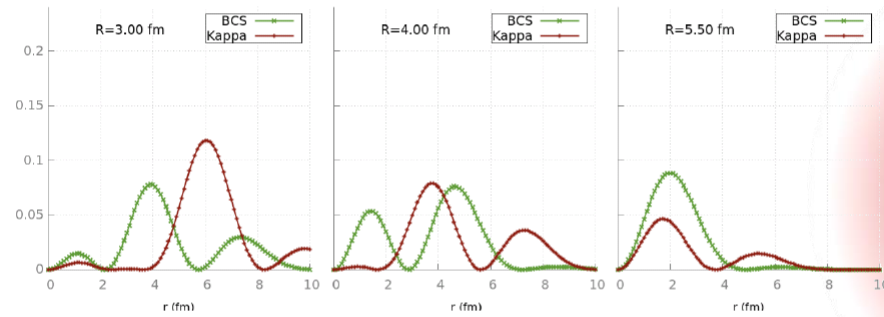
$$\Psi_N = \Psi(\mathbf{r}_1\sigma_1, \mathbf{r}_2\sigma_2, \dots, \mathbf{r}_N\sigma_N) = \mathcal{A}[\phi(\mathbf{r}_1\sigma_1; \mathbf{r}_2\sigma_2)\phi(\mathbf{r}_3\sigma_3; \mathbf{r}_4\sigma_4) \cdots \phi(\mathbf{r}_{N-1}\sigma_{N-1}; \mathbf{r}_N\sigma_N)]$$



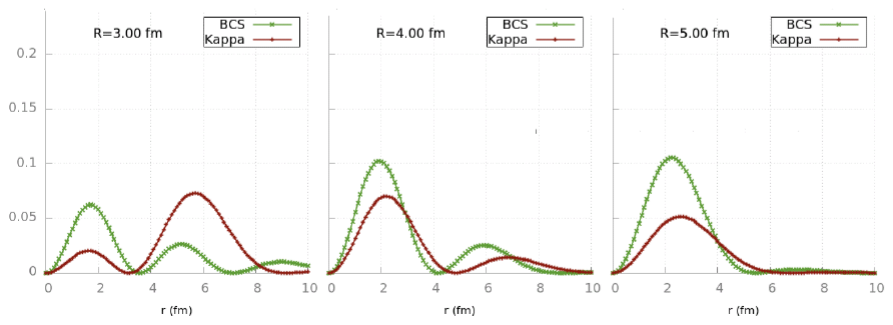
Covariant HFB with projection on good particle number prior to variation



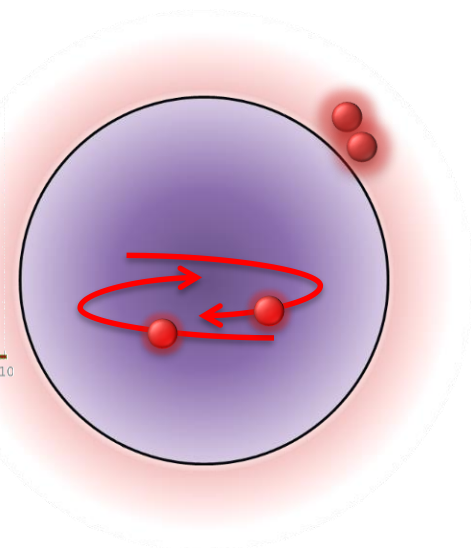
<sup>22</sup>O, N-RHB

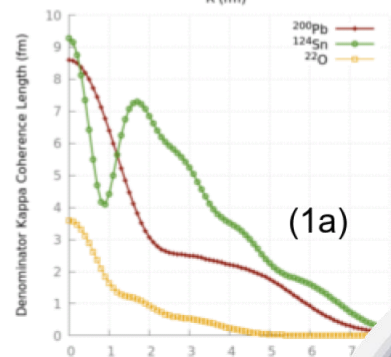
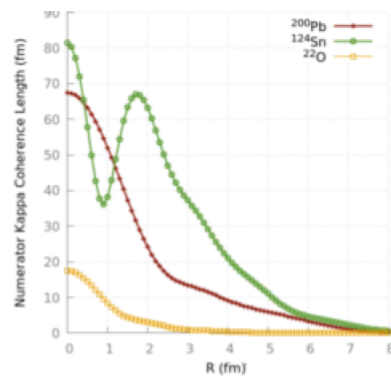
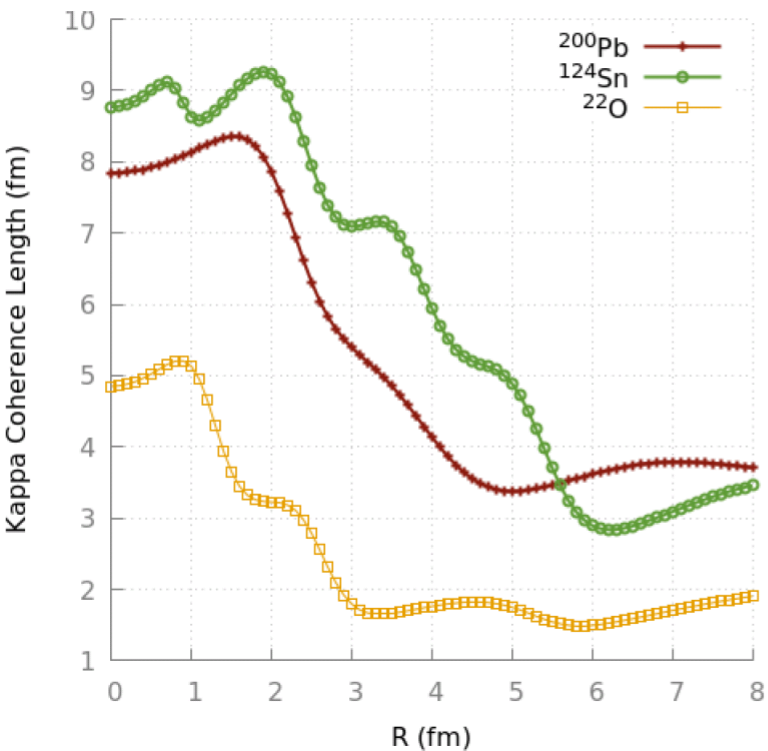


<sup>124</sup>Sn, N-RHB



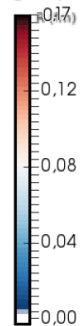
<sup>200</sup>Pb, N-RHB



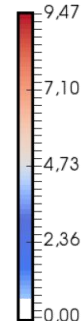


(1a)

(2a)

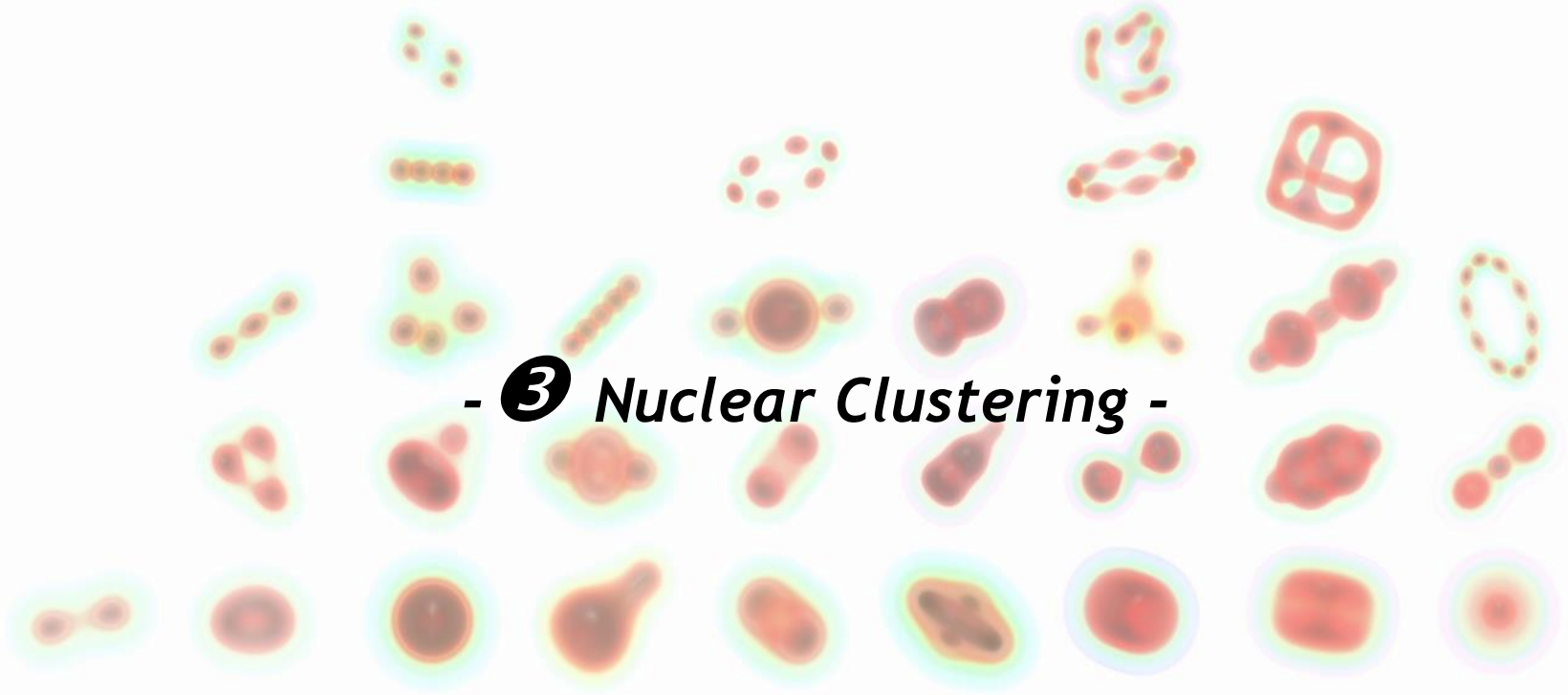


(1b)



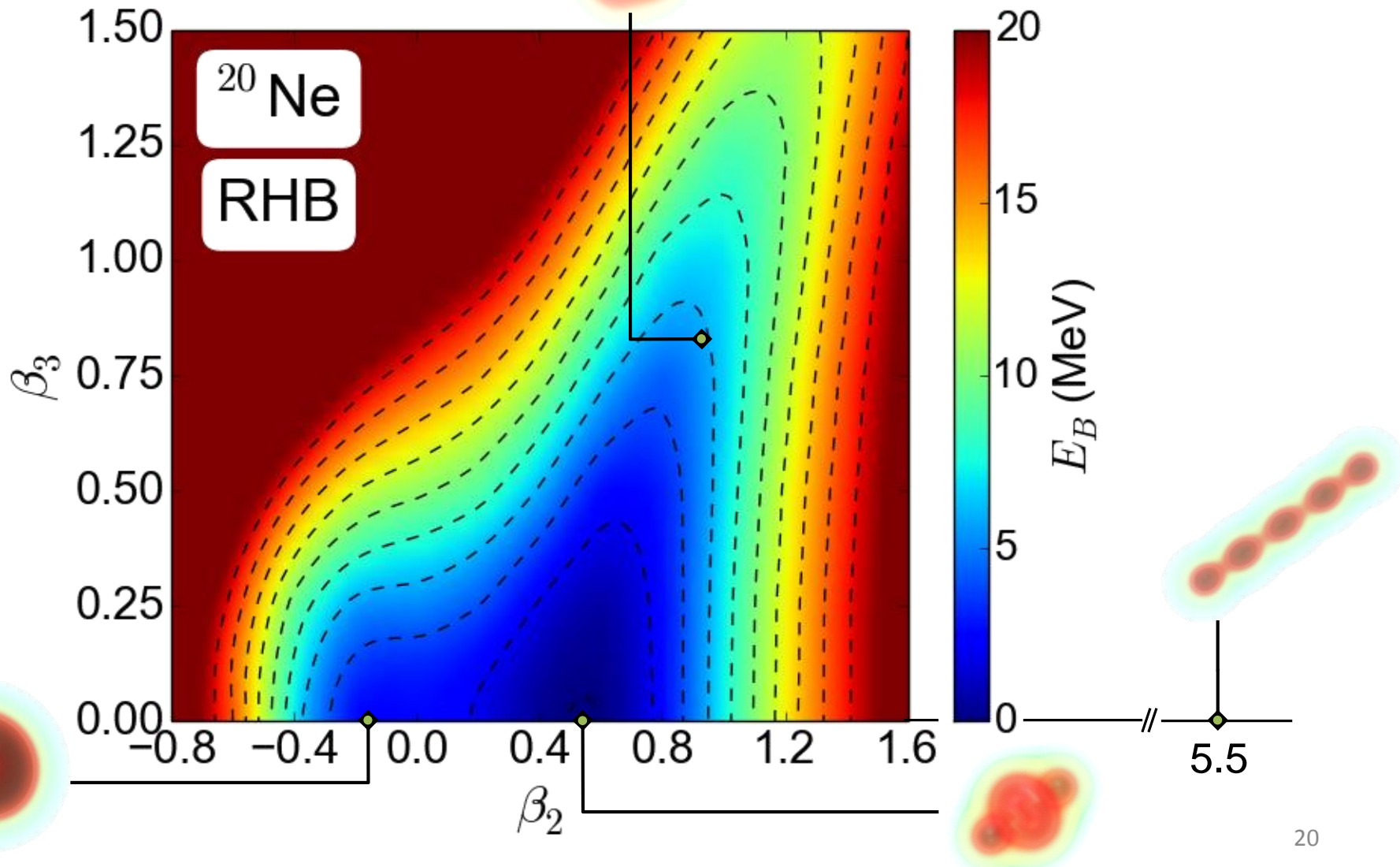
(2b)

- **3** Nuclear Clustering -



✦ First indicator: mass density at the MF level

Marevic *et al*, submitted to PRC

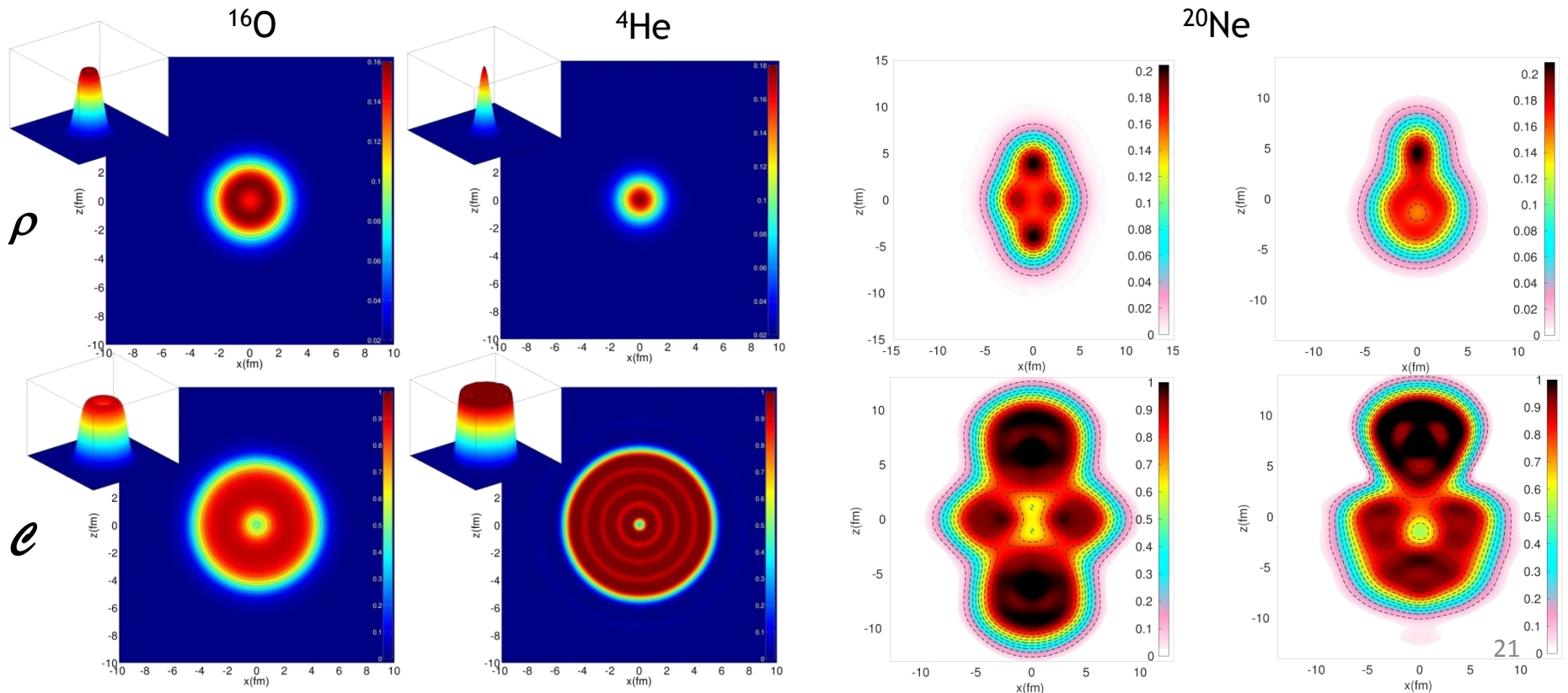


⇒ Conditional probability  $R_{q\sigma}(\mathbf{r}, \mathbf{r}') = \rho_{q\sigma}(\mathbf{r}') - \frac{|\rho_{q\sigma\sigma}(\mathbf{r}, \mathbf{r}')|^2}{\rho_{q\sigma}(\mathbf{r})}$

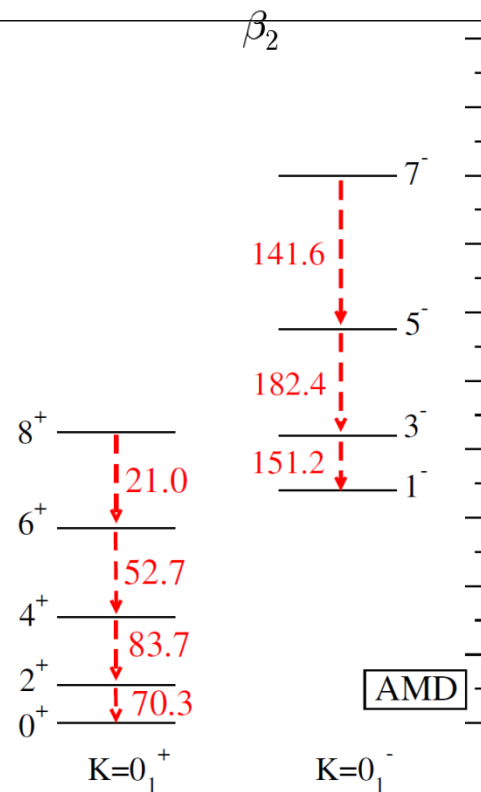
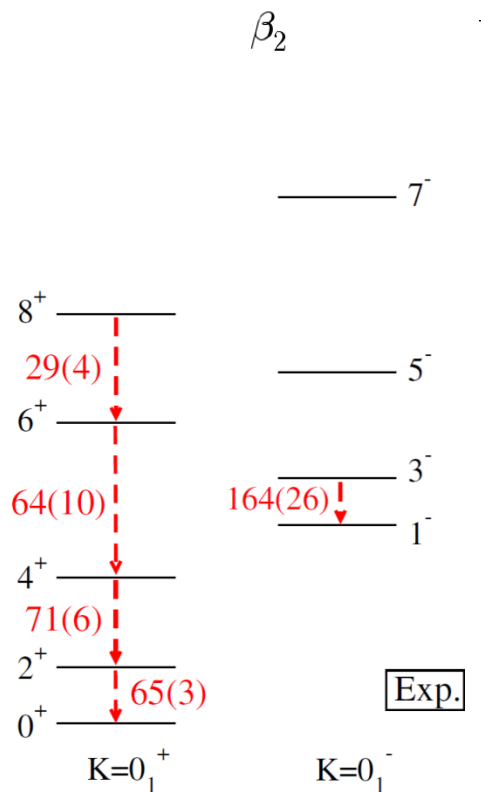
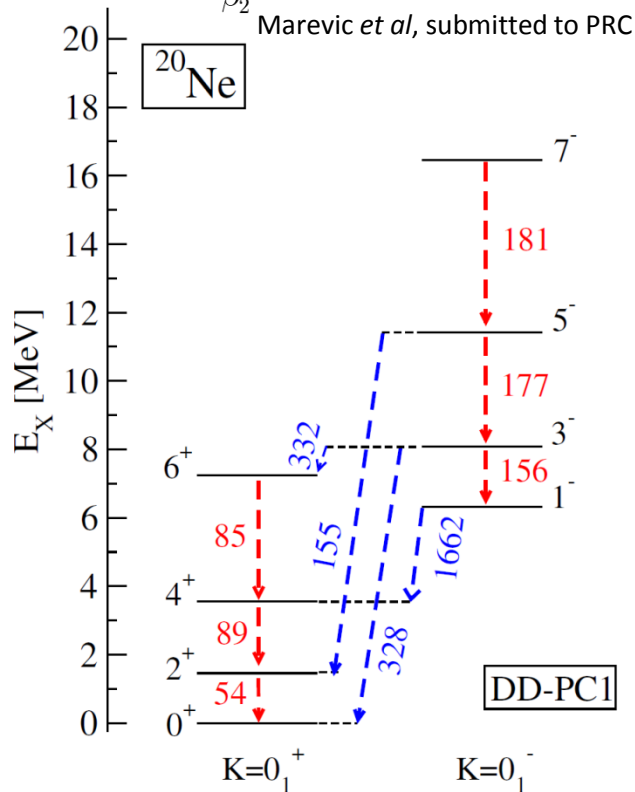
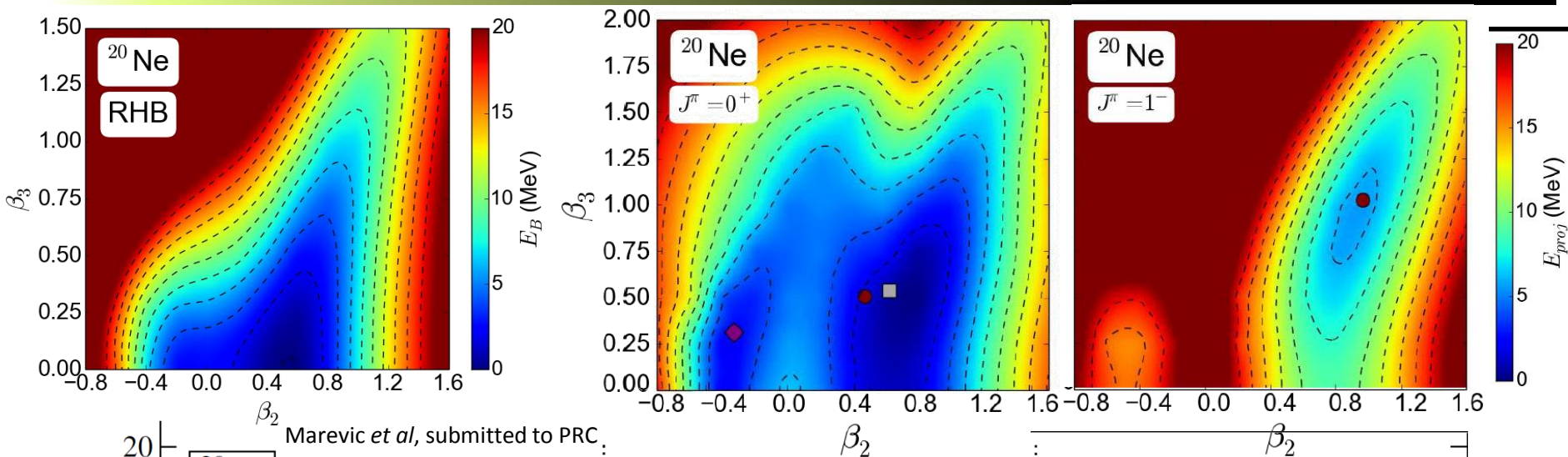
Reinhard et al, Phys. Rev. C 83, 034312 (2011)

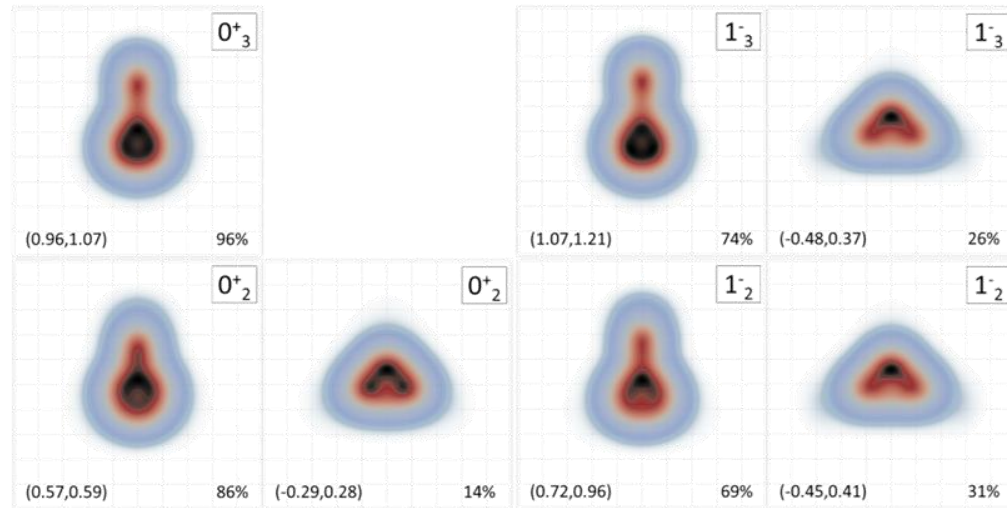
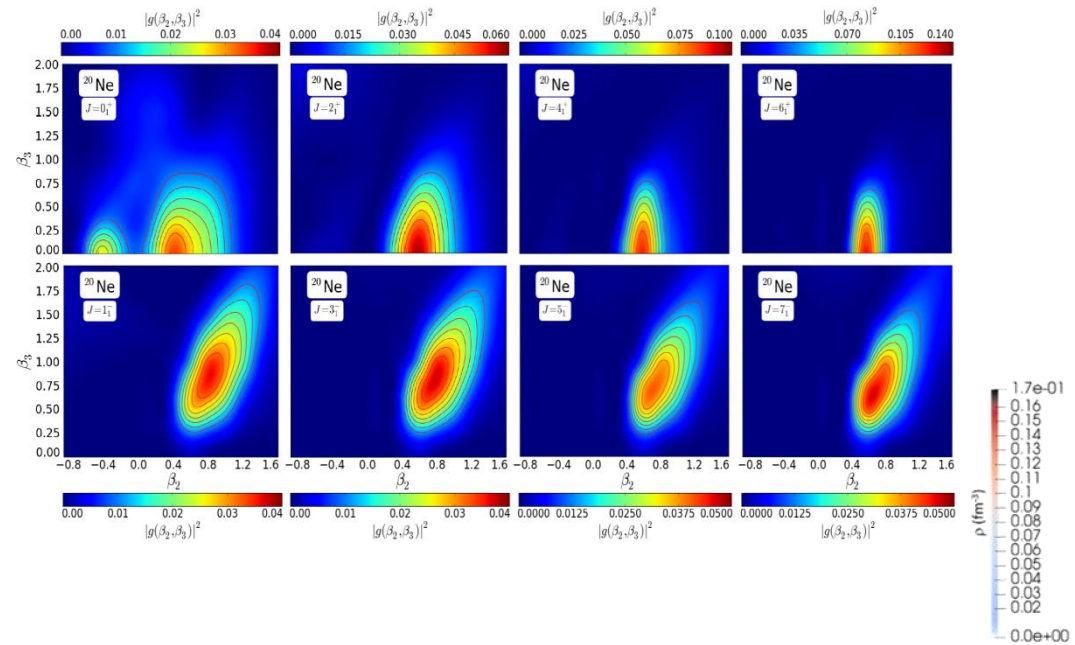
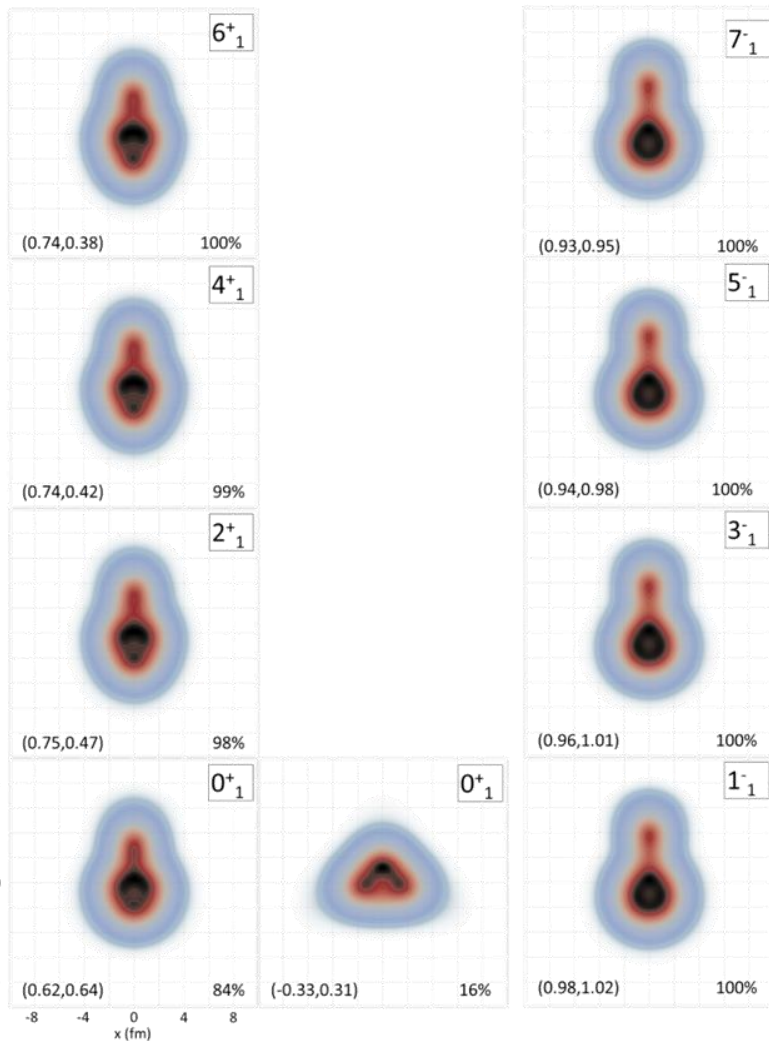
$$C_{q\sigma}(\mathbf{r}) = \left[ 1 + \left( \frac{\tau_{q\sigma} \rho_{q\sigma} - \frac{1}{4} [\nabla \rho_{q\sigma}]^2 - \mathbf{j}_{q\sigma}^2}{\rho_{q\sigma} \tau_{q\sigma}^{\text{TF}}} \right)^2 \right]^{-1}$$

0.5 : signals a nearly homogeneous Fermi gas  
 1 : localized  $\alpha$ -like state (in N=Z systems)

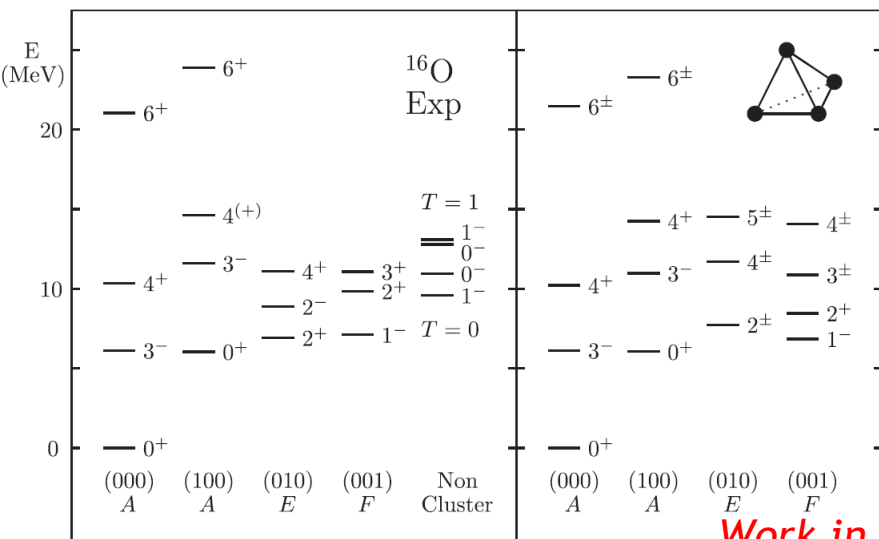
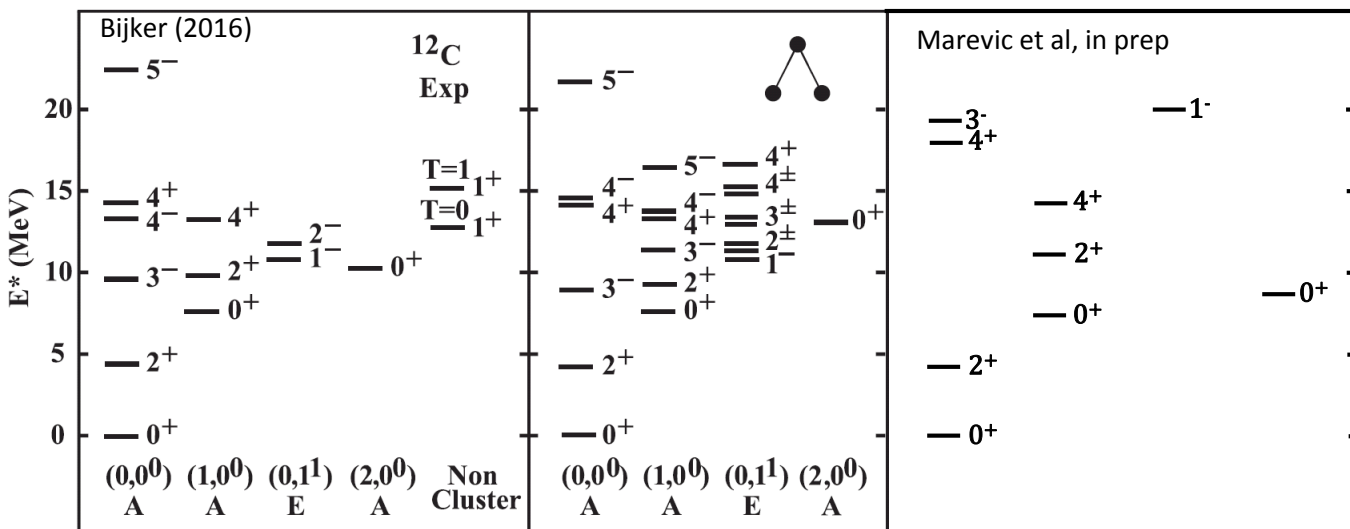


★ Beyond MF level

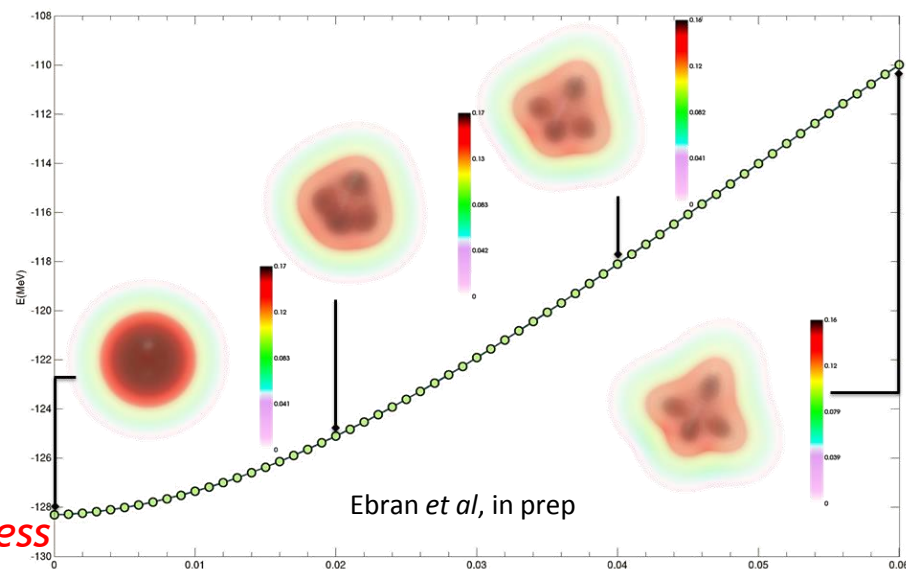




★ rotation-vibration bands



Work in progress

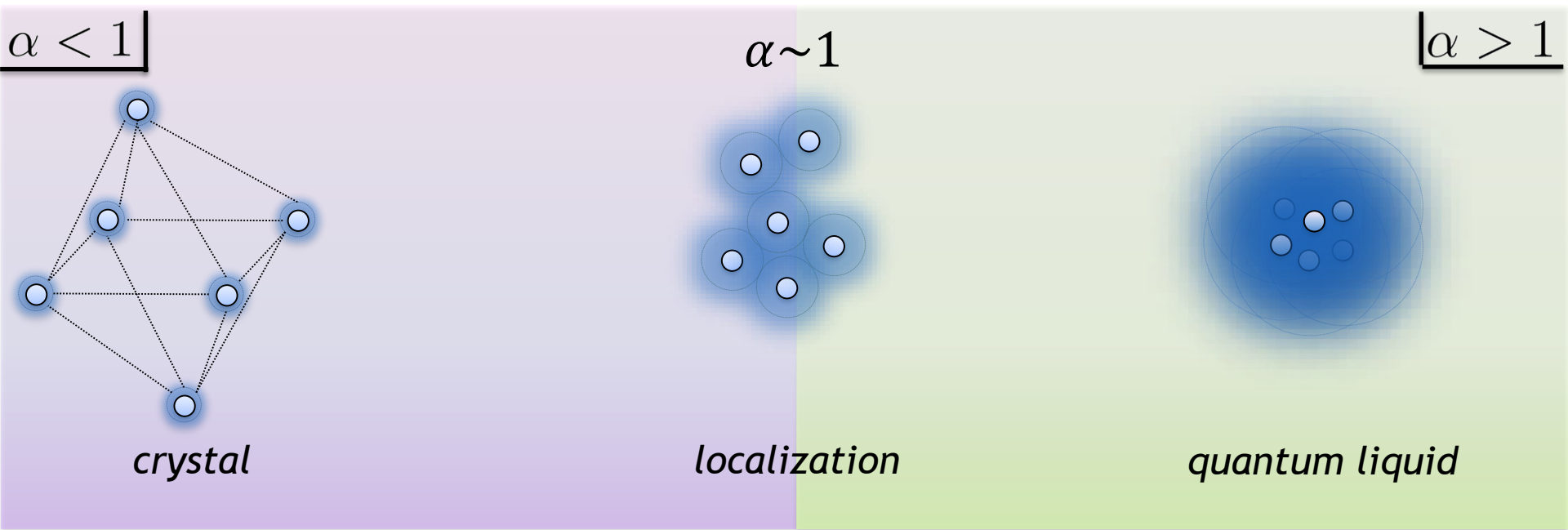
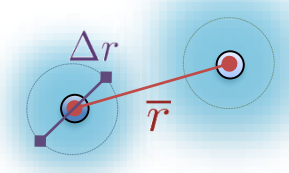




★ Guidance : the localization parameter

$$\alpha = \frac{\Delta r}{\bar{r}} = f \left( \frac{E_{kin}}{E_{pot}} \right) \propto \left[ \frac{A^{1/3}}{\bar{r} V_0^{1/2}} \right]^{\frac{1}{2}}$$

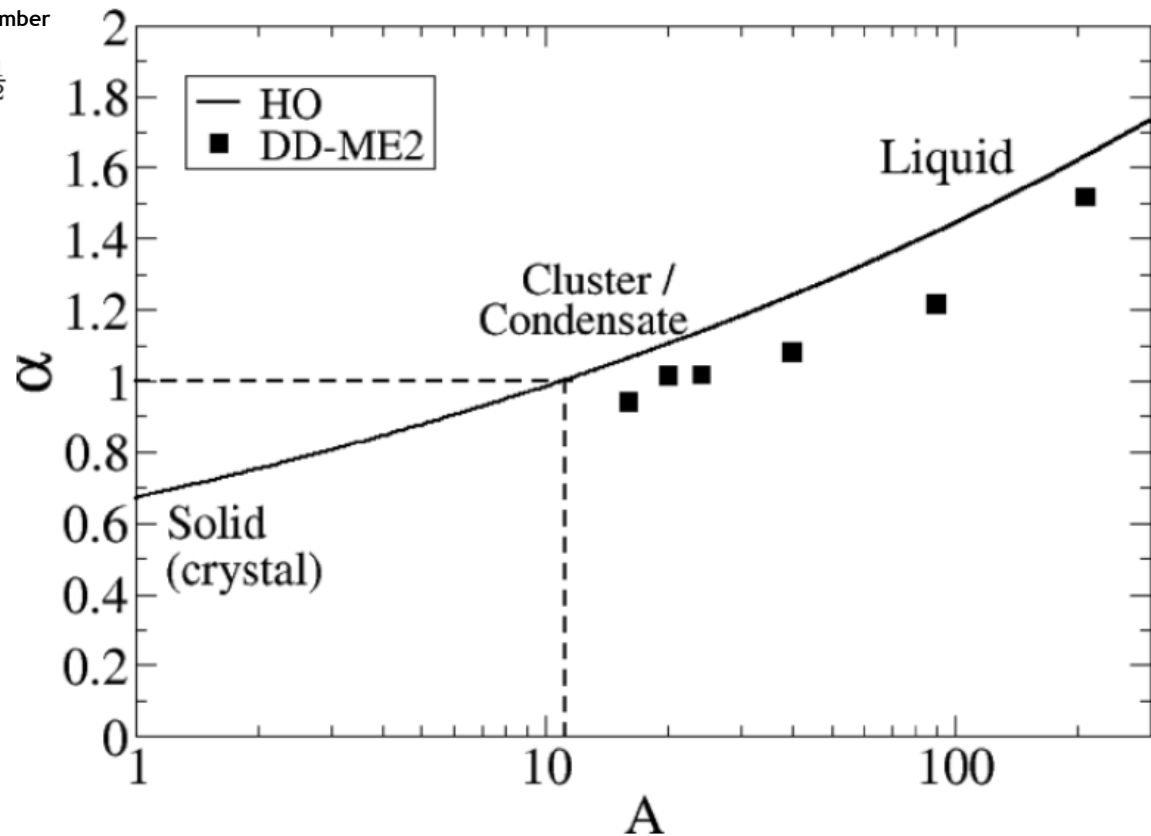
● particle number  
 ● depth of the confining potential  
 ● average inter-particle distance  $\Rightarrow$  density



✦ Influence of the particle number

$$\alpha = \frac{\Delta r}{\bar{r}} = f \left( \frac{E_{kin}}{E_{pot}} \right) \propto \left[ \frac{A^{1/3}}{\bar{r} V_0^{1/2}} \right]^{\frac{1}{2}}$$

○ particle number

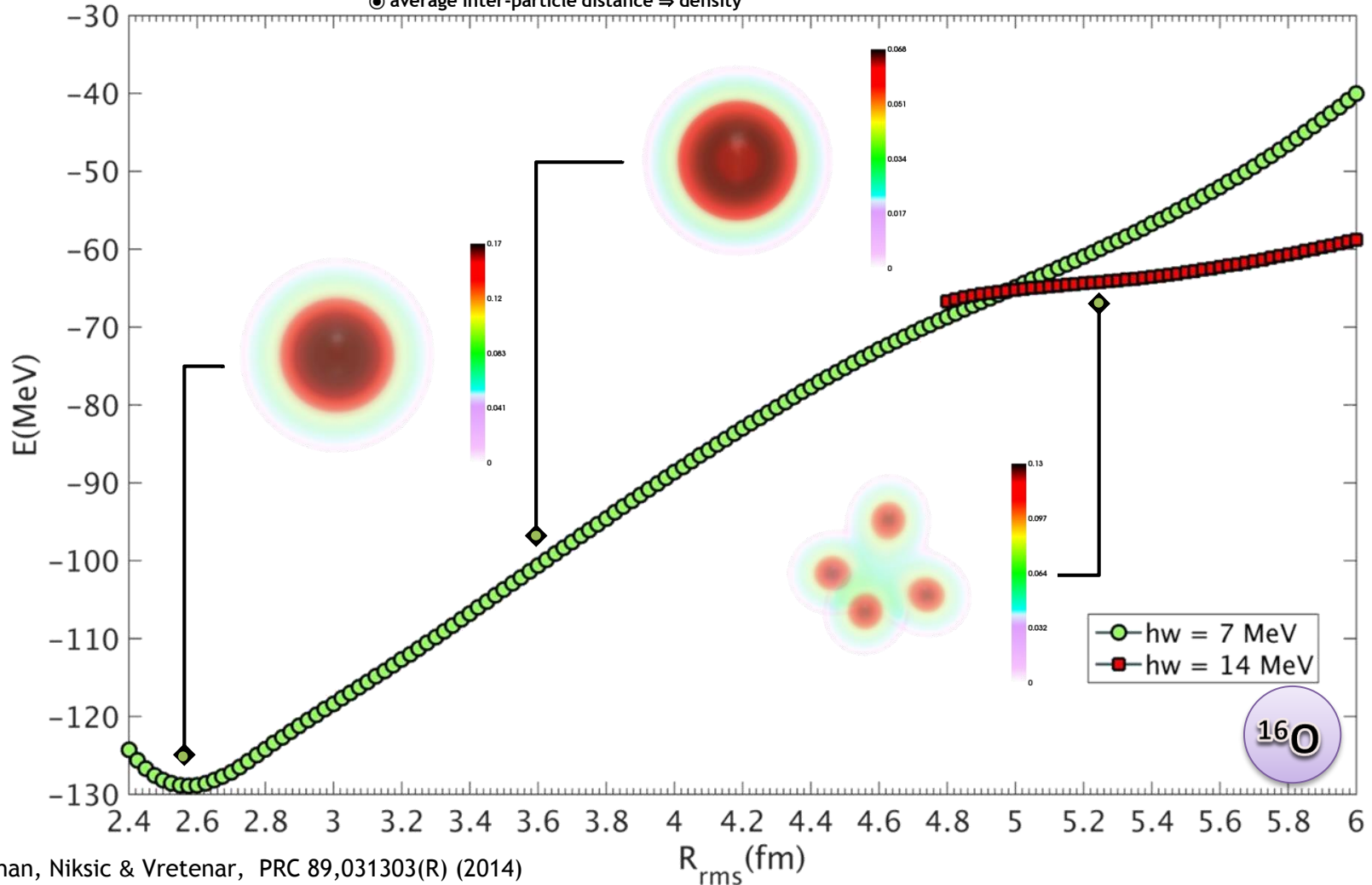


➔ Clustering is more likely to be found in light systems

★ Influence of the density

$$\alpha = \frac{\Delta r}{\bar{r}} = f \left( \frac{E_{kin}}{E_{pot}} \right) \propto \left[ \frac{A^{1/3}}{\bar{r} V_0^{1/2}} \right]^{\frac{1}{2}}$$

○ average inter-particle distance  $\Rightarrow$  density



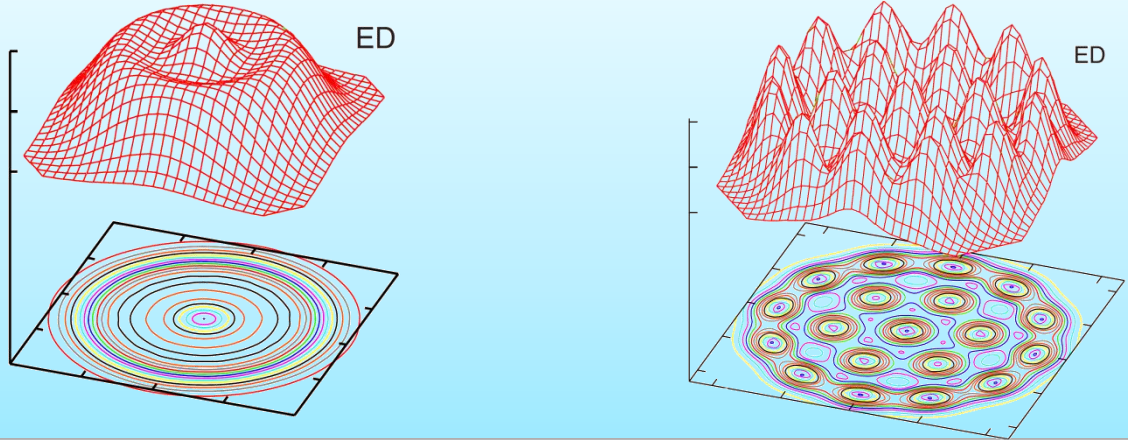
⊛ Influence of the confining potential

$$\alpha = \frac{\Delta r}{\bar{r}} = f\left(\frac{E_{kin}}{E_{pot}}\right) \propto \left[\frac{A^{1/3}}{\bar{r}V_0^{1/2}}\right]^{\frac{1}{2}}$$

⊙ depth of the confining potential



⇒ Electrons in quantum dots

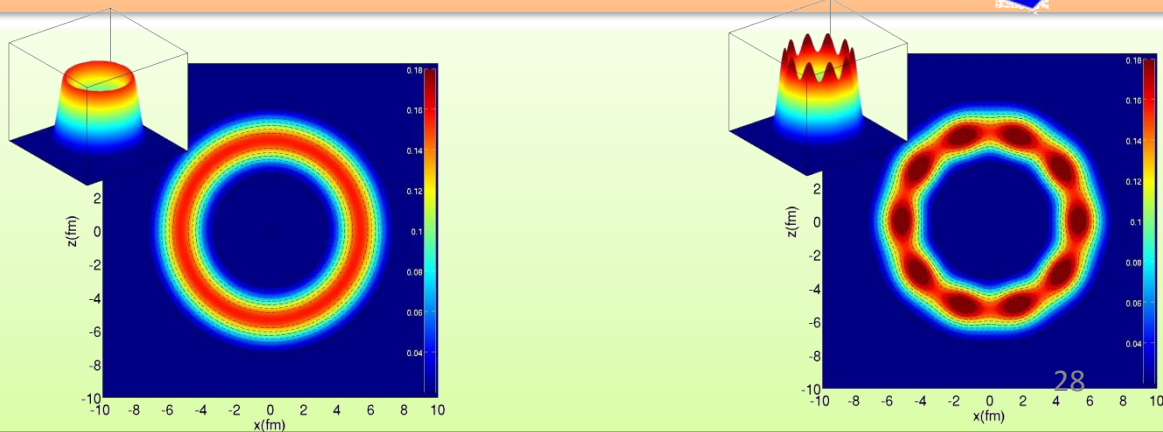


Yannouleas and Landman, Rep.Prog.Phys. 70, 2067 (2007)

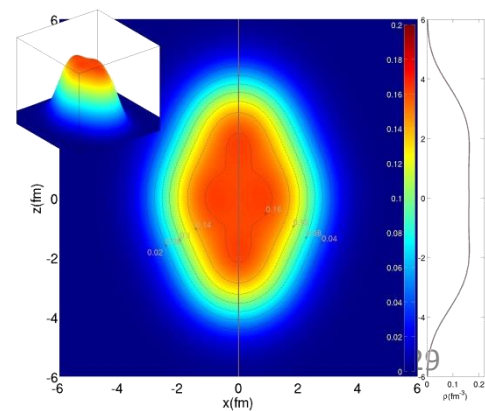
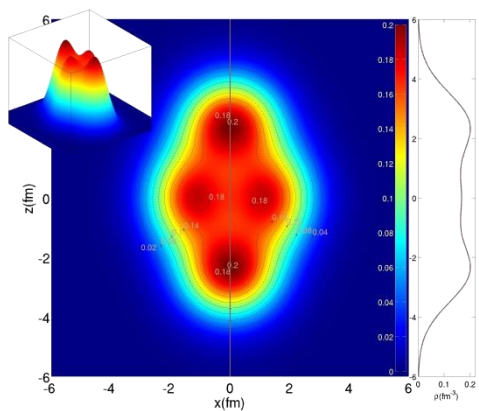
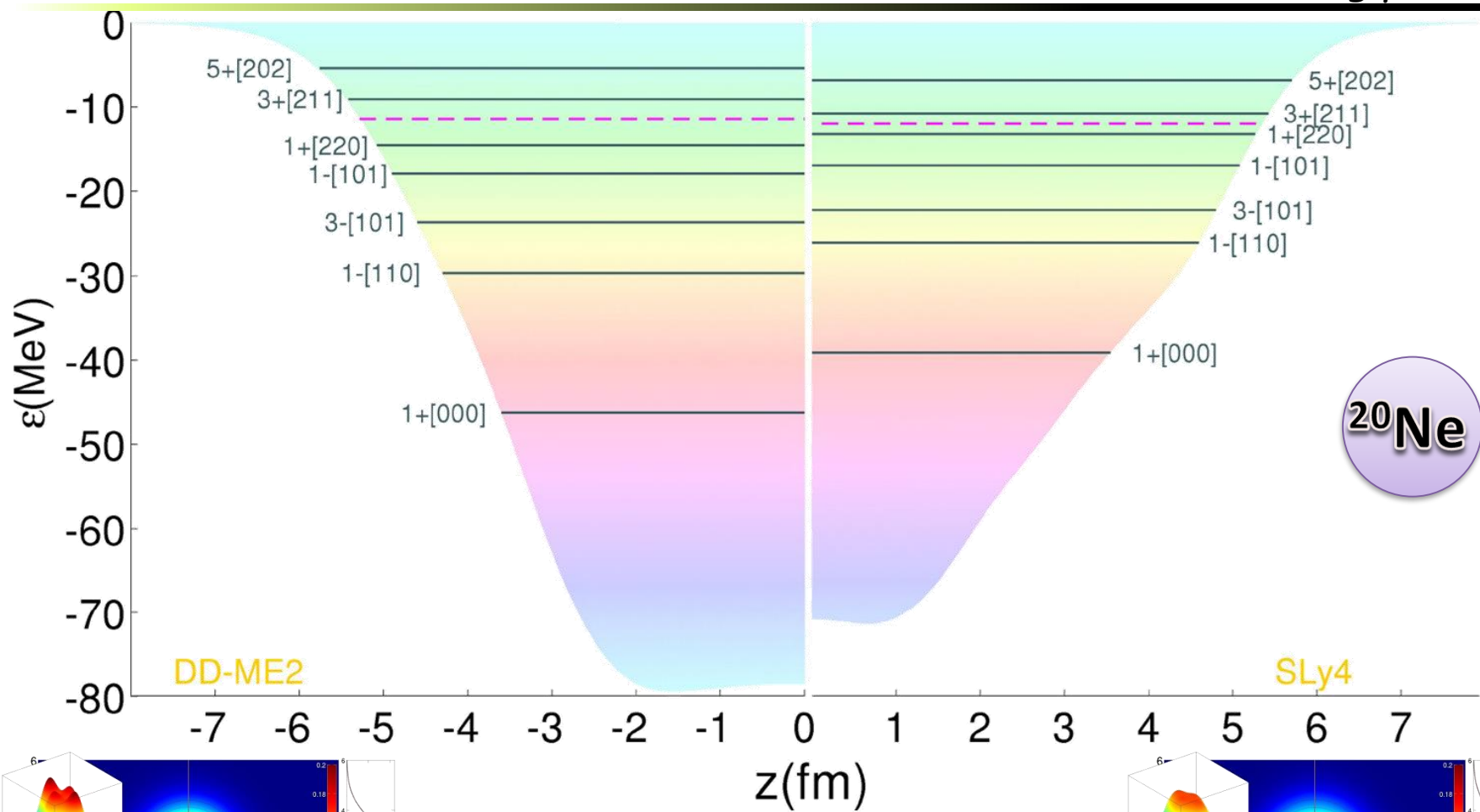
⇒ Neutral bosons in rotating trap



⇒ Nucleons in <sup>40</sup>Ca



⊛ Influence of the confining potential

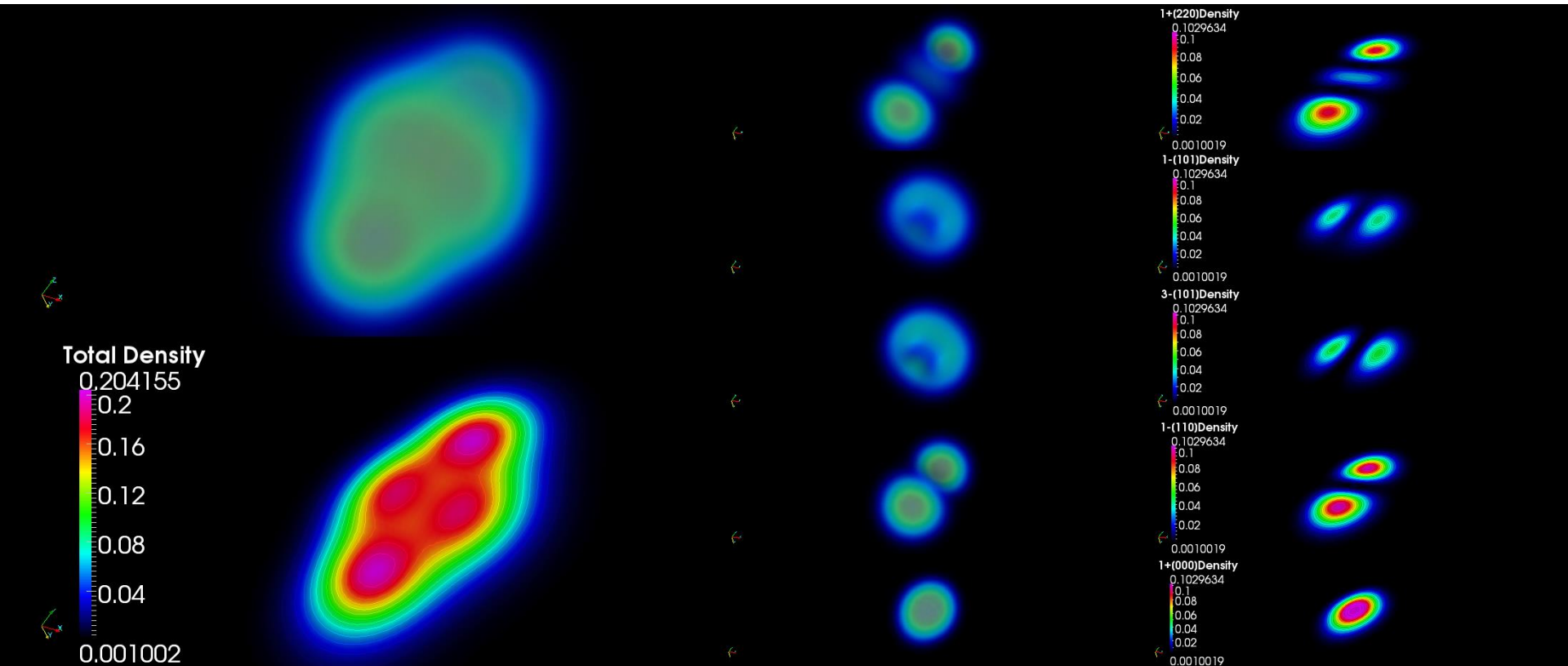


✦ Influence of the confining potential



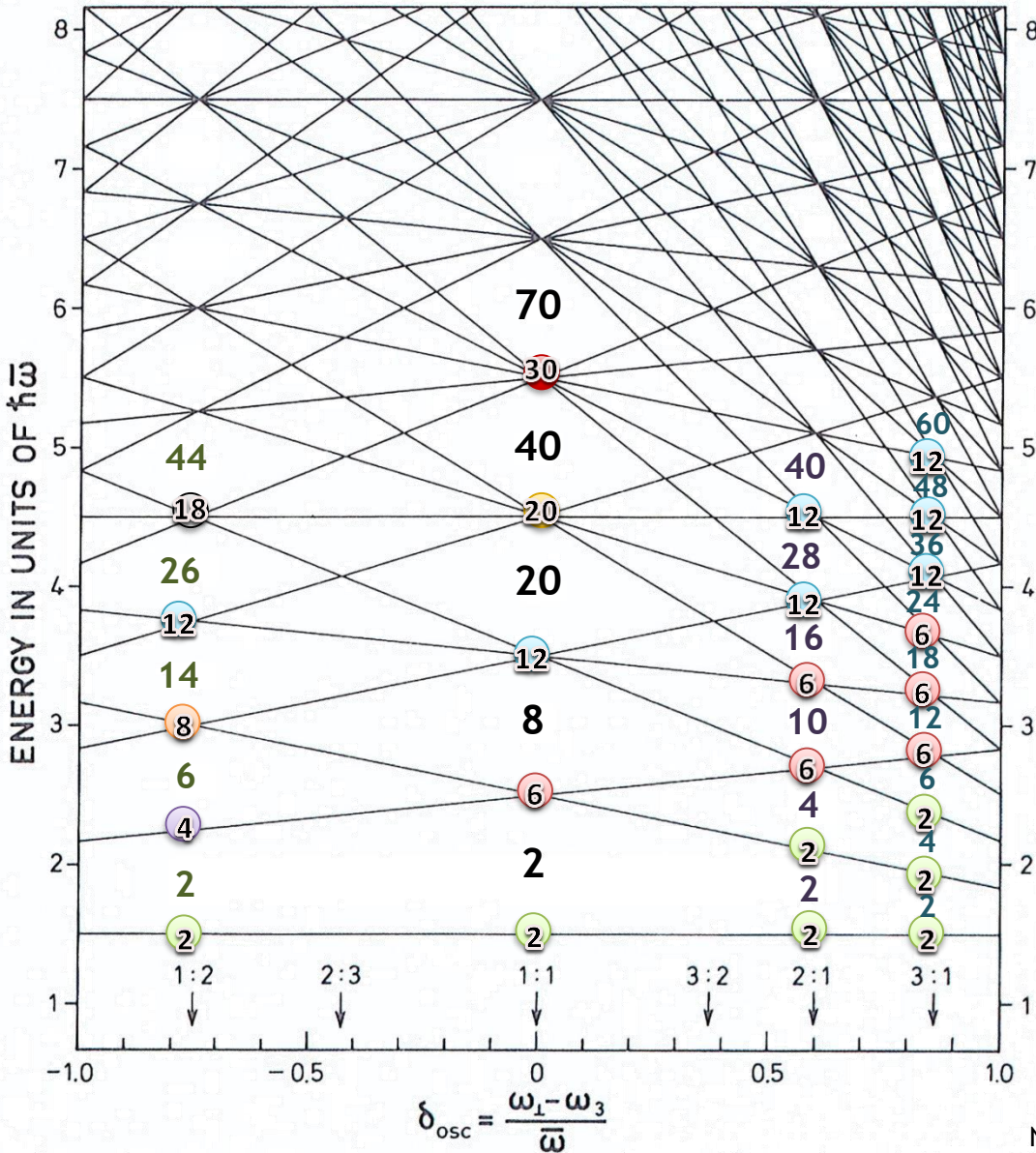
$$\alpha = \frac{\Delta r}{\bar{r}} = f \left( \frac{E_{kin}}{E_{pot}} \right) \propto \left[ \frac{A^{1/3}}{\bar{r} V_0^{1/2}} \right]^{\frac{1}{2}}$$

● depth of the confining potential

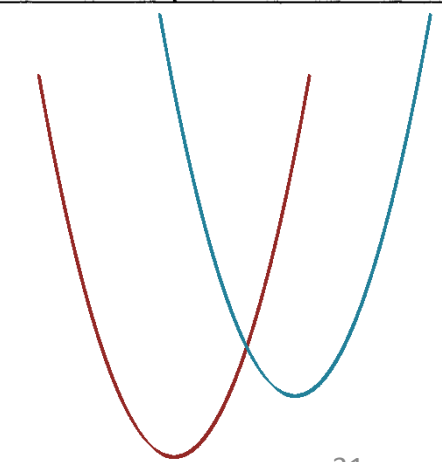


⊛ Influence of the confining potential

Harmonic oscillator case

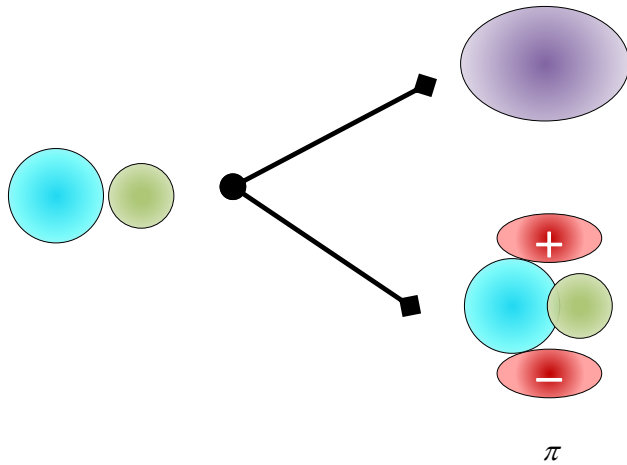


SPHERICAL MAGIC NUMBERS	SUPERDEFORMED PROLATE MAGIC NUMBERS	SUPERDEFORMED PROLATE SPECTRUM
70 ○	→ ○○ 140	4 —
40 ○	→ ○○ 110	4 —
20 ○	→ ○○ 80	3 —
8 ○	→ ○○ 60	3 —
2 ○	→ ○○ 40	2 —
	→ ○○ 28	2 —
	→ ○○ 16	1 —
	→ ○○ 10	1 —
	→ ○○ 4	0 —
	→ ○ 2	0 —
	<i>A</i> <i>B</i>	(000) (001)



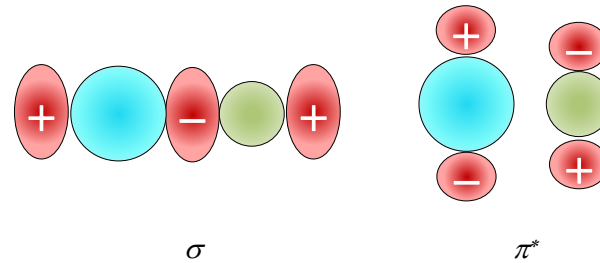
⊛ Influence of the neutron excess

2-center cluster in a N=Z system



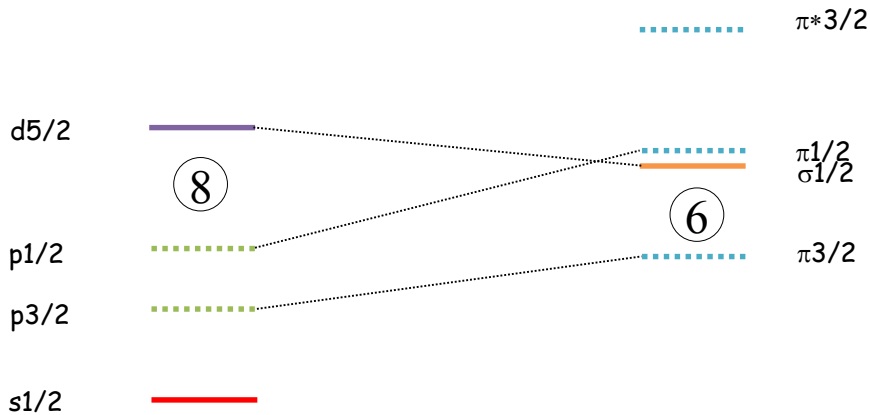
Neutron rich isotope

melting



Covalent bonding

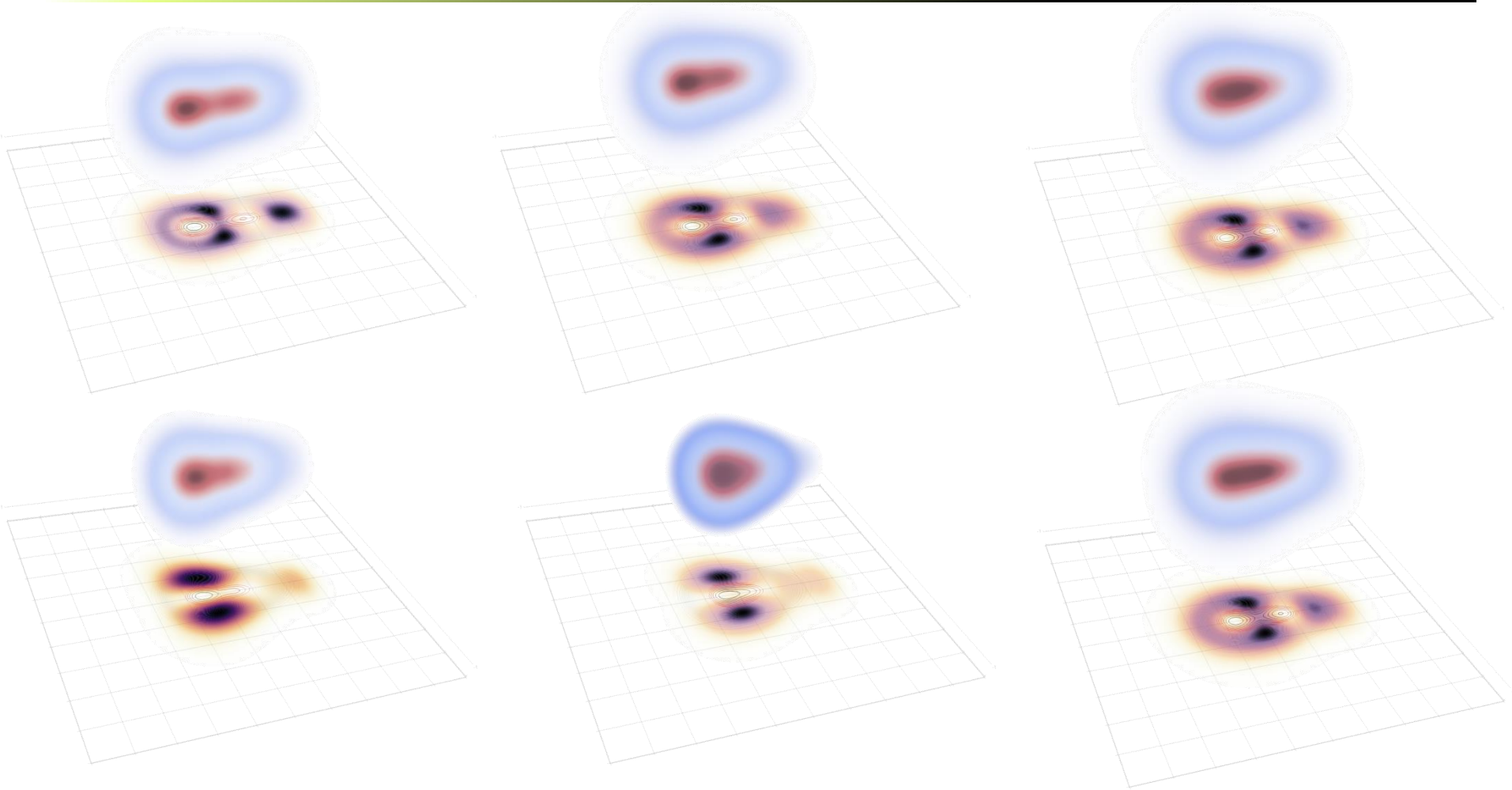
Be case



N=8 magic number breaking



⊛ Influence of the neutron excess



$0^+_2$

$0^+_1$



$^{10}\text{Be}$

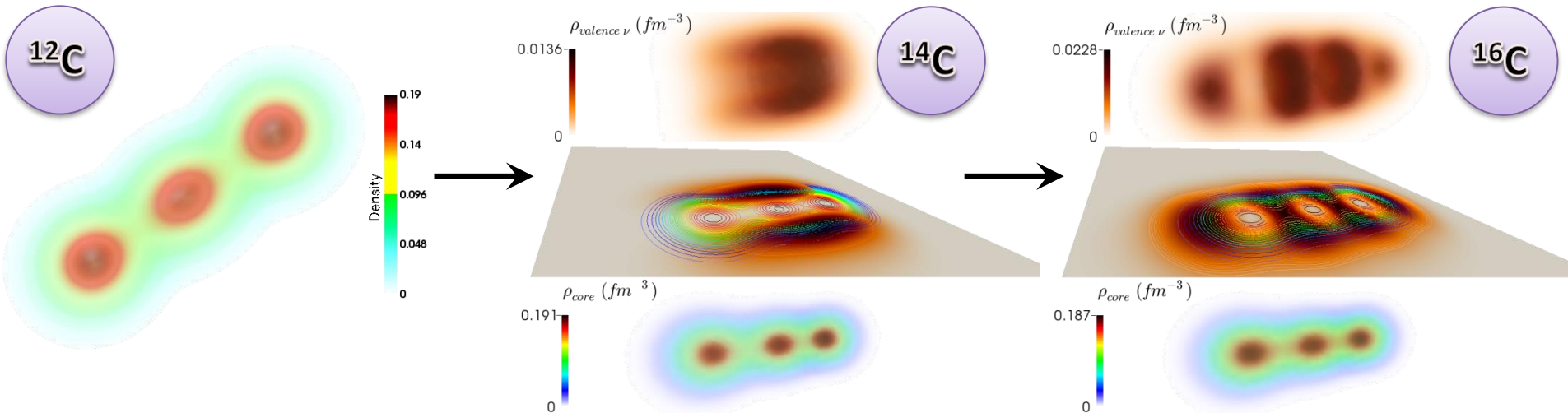


$^{12}\text{Be}$

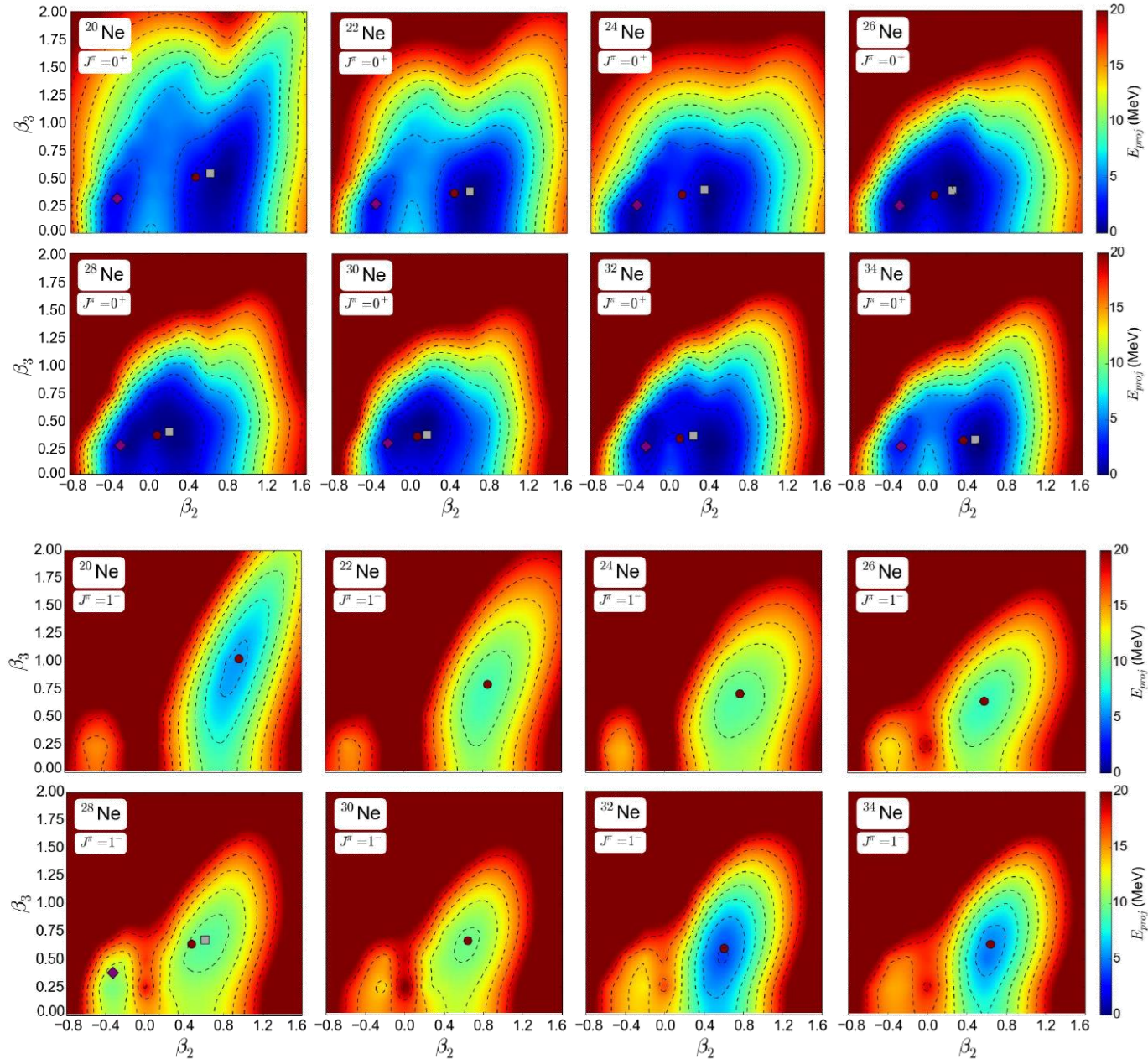


$^{14}\text{Be}$

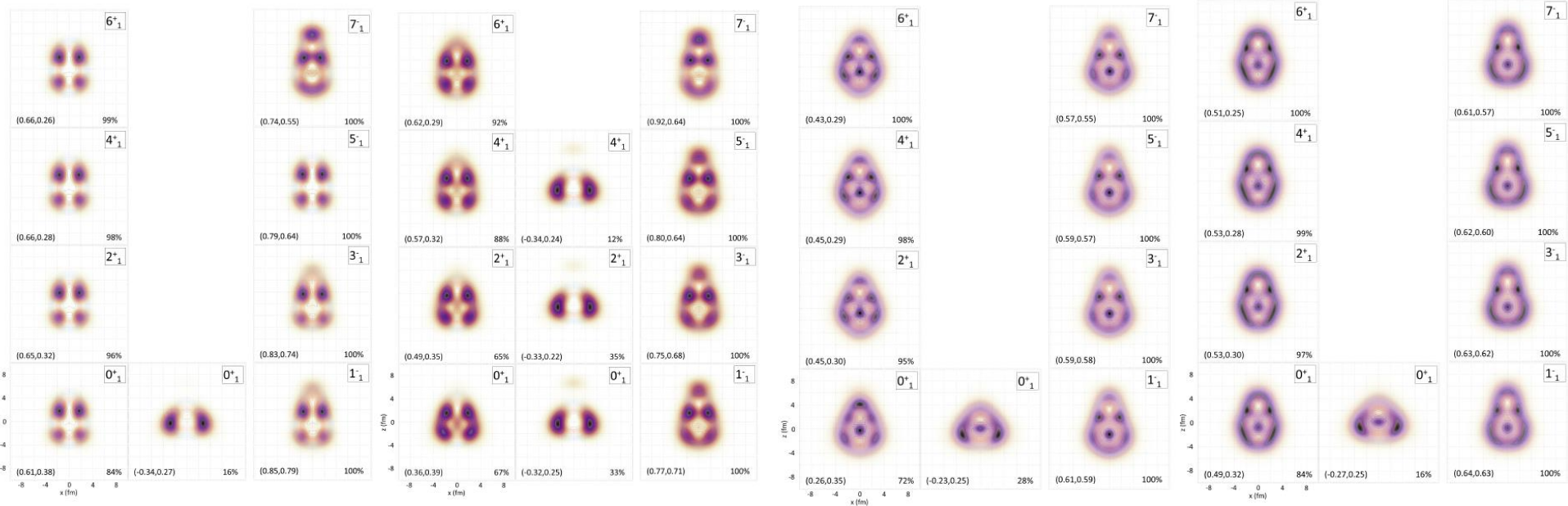
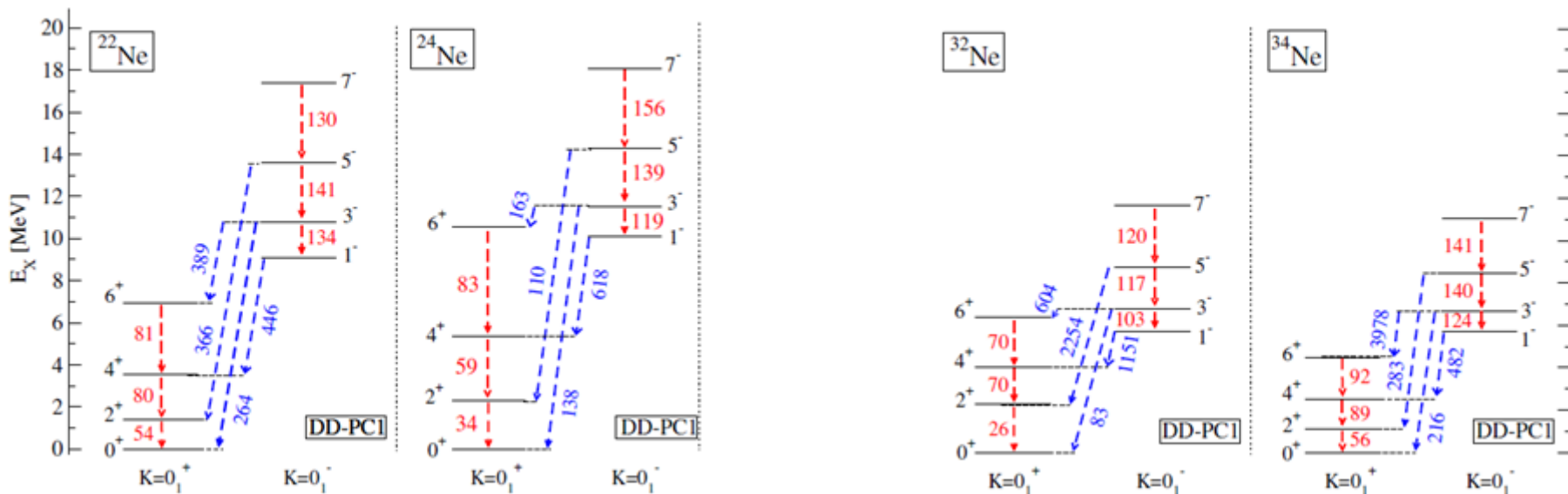
⊛ Influence of the neutron excess



⊛ Influence of the neutron excess



⊛ Influence of the neutron excess

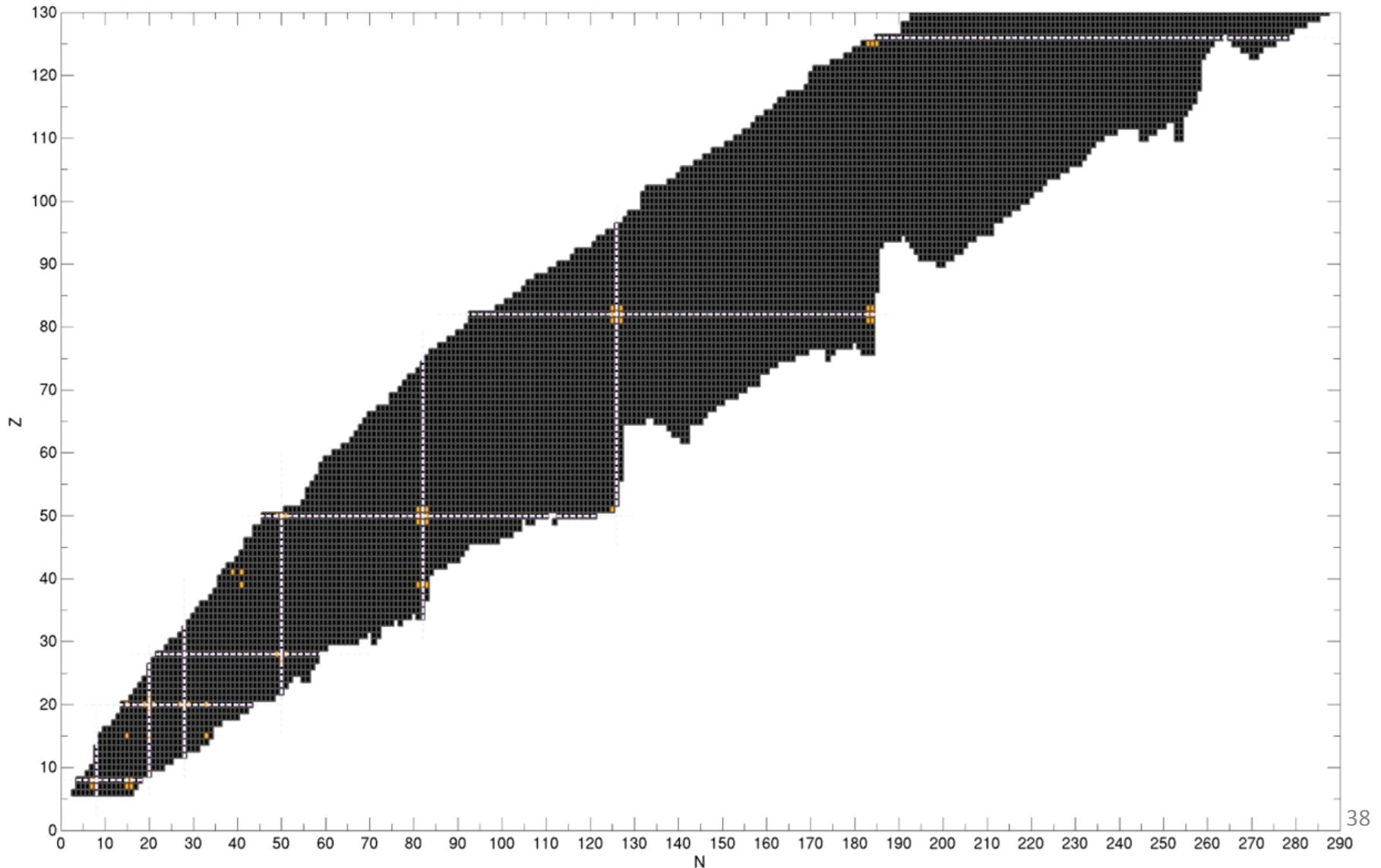


# Conclusion & Perspectives

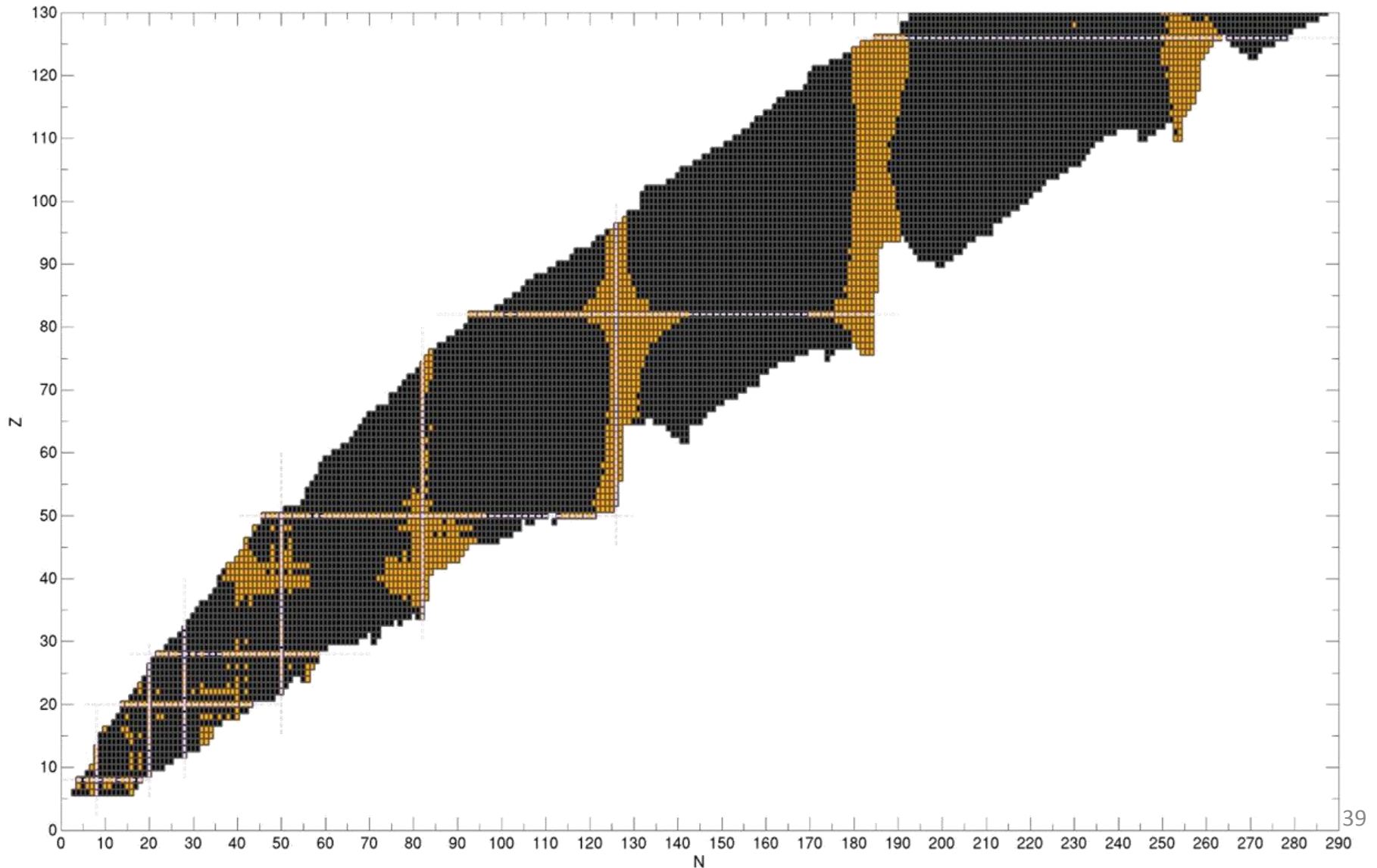
- ⇒ Nuclear EDFs frame the various nuclear properties in a unified and consistent way
- ⇒ Key feature : breaking/restoration of symmetries to efficiently account for nondynamical correlations
- ⇒ Di-neutron like configuration at the surface of superfluid nuclei
- ⇒ Clustering in ground and excited states of nuclei : impact of the average density and the depth of the confining potential
- ⇒ Covalent bonding in neutron rich systems
- ⇒ Particle number projection, conditional probability distribution, form factor, quarteting ... under development

Thank you for your attention

⇒ Good description of the ground state of ~50 nuclei



Breaking  $U(1)$  : good description of the ground state of  $\sim 300$  nuclei



⇒ Breaking  $U(1)$  and  $O(3)$  : good description of the ground state of all nuclei

